

# On the Measurement of Elemental Abundance Ratios in Inner Galaxy H II Regions

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## ABSTRACT

Although abundance gradients in the Milky Way Galaxy certainly exist, details remain uncertain, particularly in the inner Galaxy, where stars and H II regions in the Galactic plane are obscured optically. In this paper we revisit two previously studied, inner Galaxy H II regions: G333.6–0.2 and W43. We observed three new positions in G333.6–0.2 with the Kuiper Airborne Observatory and reobserved the central position with the Infrared Space Observatory's Long Wavelength Spectrometer in far-infrared lines of S<sup>++</sup>, N<sup>++</sup>, N<sup>+</sup>, and O<sup>++</sup>. We also added the N<sup>+</sup> lines at 122 and 205  $\mu\text{m}$  to the suite of lines measured in W43 by Simpson et al. (1995). The measured electron densities range from  $\sim 40$  to over  $4000\text{ cm}^{-3}$  in a single H II region, indicating that abundance analyses must consider density variations, since the critical densities of the observed lines range from 40 to  $9000\text{ cm}^{-3}$ . We propose a method to handle density variations and make new estimates of the S/H and N/H abundance ratios. We find that our sulfur abundance estimates for G333.6–0.2 and W43 agree with the S/H abundance ratios expected for the gradient previously reported by Simpson et al., with the S/H values revised to be smaller owing to changes in collisional excitation cross sections. The estimated N/H, S/H, and N/S ratios are the most reliable because of their small corrections for unseen ionization states ( $\lesssim 10\%$ ). The estimated N/S ratios for the two sources are smaller than what would be calculated from the N/H and S/H ratios in our previous paper. If all low excitation H II regions had

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similar changes to their N/S ratios as a result of adding measurements of  $N^+$  to previous measurements of  $N^{++}$ , there would be no or only a very small gradient in N/S. This is interesting because nitrogen is considered to be a secondary element and sulfur is a primary element in galactic chemical evolution calculations. We compute models of the two H II regions to estimate corrections for the other unseen ionization states. We find, with large uncertainties, that oxygen does not have a high abundance, with the result that the N/O ratio is as high ( $\sim 0.35$ ) as previously reported. The reasons for the uncertainty in the ionization corrections for oxygen are both the non-uniqueness of the H II region models and the sensitivity of these models to different input atomic data and stellar atmosphere models. We discuss these predictions and conclude that only a few of the latest models adequately reproduce H II region observations, including the well-known, relatively-large observed  $Ne^{++}/O^{++}$  ratios in low- and moderate-excitation H II regions.

*Subject headings:* Galaxy: abundances — H II regions — ISM: individual (G333.6–0.2, W43)

## 1. Introduction

The presence of radial abundance gradients in the plane of the Milky Way is now an established fact and is seen in both stars and gaseous nebulae (see, e.g., Henry & Worthey 1999; Rolleston et al. 2000; and references in both papers). The cause is generally thought to be due to the formation history and subsequent evolution of the Galaxy; thus the observed gradients are a major tool for understanding this history (see, e.g., Hou, Prantzos, & Boissier 2000; Chiappini, Matteucci, & Romano 2001; Chiappini, Romano, & Matteucci 2003, and references therein). The basic idea is that the inner Galaxy formed before the outer Galaxy and the higher molecular gas density in the inner Galaxy produces a higher star formation rate. The result is a greater return to the interstellar medium (ISM) in the inner Galaxy of both “primary” alpha elements from massive star supernovae and “secondary” elements like nitrogen. Secondary nitrogen is produced by CNO burning of already existing carbon and oxygen in intermediate-mass stars and subsequently returned to the ISM through mass loss. However, both the wide range of uncertain input parameters to chemical evolution models (Pagel 2001) and the uncertain details of the abundance variation of each element, primary and secondary, contribute to our less-than-complete understanding of the formation and evolution of the Milky Way. In this paper we try to improve our knowledge of the abundances in inner Galaxy H II regions and also test the methods that have been used to

date to determine abundance gradients from low-excitation, optically obscured H II regions.

H II regions in the inner Galaxy (galactocentric radius  $R_G \lesssim 6$  kpc, where  $R_\odot = 8$  kpc) are not accessible to optical observers because of interstellar extinction. Consequently, the abundances must be determined from far-infrared (FIR) lines, such as those of  $N^+$  at 122 and 205  $\mu\text{m}$ ,  $N^{++}$  at 57  $\mu\text{m}$ ,  $O^{++}$  at 52 and 88  $\mu\text{m}$ ,  $Ne^+$  at 12.8  $\mu\text{m}$ ,  $Ne^{++}$  at 15.5 and 36  $\mu\text{m}$ ,  $S^{++}$  at 18.7 and 33.5  $\mu\text{m}$ , and  $S^{3+}$  at 10.5  $\mu\text{m}$ . Large surveys include those of Simpson et al. (1995a, hereafter SCREH) and Afflerbach et al. (1997), who observed the FIR lines (wavelengths  $> 17$   $\mu\text{m}$ ) with the Kuiper Airborne Observatory (KAO) in 18 inner Galaxy H II regions, and Martín-Hernández et al. (2002a), who observed 13 H II regions with the Infrared Space Observatory (ISO). The KAO observations included FIR measurements of the lines with wavelength  $> 18$   $\mu\text{m}$ , whereas the ISO observations included the mid-infrared lines, measured with ISO's Short Wavelength Spectrometer, SWS, but not the [N II] 205  $\mu\text{m}$  line.

A major result of these studies is that inner Galaxy H II regions have low excitation compared to typical H II regions at larger  $R_G$ ; this is apparent in both heavy element ionic ratios such as  $O^{++}/S^{++}$  (SCREH) and  $Ne^{++}/Ne^+$  (Simpson & Rubin 1990; Giveon et al. 2002), and also in  $He^+/H^+$  measured from radio recombination lines (Churchwell et al. 1978; Thum, Mezger, & Pankonin 1980). It is still controversial whether the observed increase in excitation with increasing  $R_G$  is due entirely to heavy element opacity effects in the H II regions and stellar atmospheres, or whether there is an additional gradient in the maximum stellar effective temperature,  $T_{\text{eff}}$ , of the exciting stars as well (e.g., Giveon et al. 2002; Martín-Hernández et al. 2002b; Smith, Norris, & Crowther 2002; Morisset et al. 2003).

Two of the abundance gradients that are most needed for galactic chemical evolution studies are the O/H and N/H ratios, representing primary and secondary elements, respectively. In optically obscured H II regions the only ionization state of oxygen available for study is  $O^{++}$ , which means that corrections are required for the unseen  $O^+$  ions. The situation is better for nitrogen, as there are [N II] lines as well as [N III] lines in the FIR wavelength range. Unfortunately, the [N II] 122  $\mu\text{m}$  line is difficult to measure for airborne observations because of atmospheric absorption, and the [N II] 205  $\mu\text{m}$  line is difficult because of poor detector sensitivity. In addition, the [N II] lines are difficult to interpret because they are collisionally de-excited at much lower densities than the lines from the other ions. Consequently, the papers to date have used only the [N III] and [O III] lines to estimate N/O and the secondary/primary abundance ratio in optically obscured H II regions (SCREH; Afflerbach et al. 1997; Martín-Hernández et al. 2002a). This is particularly troubling since a majority of the nitrogen and an even larger fraction of the oxygen are singly ionized in the high metallicity H II regions of the inner Galaxy (note that in the outer Galaxy, the lower

metallicity H II regions can be much more highly excited and thus the  $N^{++}/O^{++}$  ratio is often a good approximation to  $N/O$ , Rudolph et al. 1997). On the other hand, a majority of the primary element sulfur is doubly ionized in H II regions. Here we devise a method of analyzing the density-sensitive [N II] lines to more accurately estimate nitrogen abundances with respect to sulfur without the need for large corrections for unseen ionization states. Using the observed  $N^{++}/N^+$  ratio and calculated H II region models, we then estimate the ionization correction factors (ICFs) for oxygen.

Two of the most luminous H II regions in the Galaxy are G333.6–0.2 at  $R_G \sim 5.5$  kpc (Wilson et al. 1970) or 5.6 kpc (Caswell & Haynes 1987) and W43 (G30.8–0.0) at  $R_G \sim 4.2$  kpc (Reifenstein et al. 1970) for  $R_\odot = 8.0$  kpc. The radio fluxes from these objects require ionizing stellar luminosities of  $10^{50}$  and  $2.9 \times 10^{50}$  photons  $s^{-1}$ , respectively (see Rubin 1968b). Both H II regions are located in the Scutum-Crux spiral arm (see Vallee 2002) in the Milky Way molecular ring. Both are excited by clusters of massive stars, observed at near-infrared (NIR) wavelengths. The cluster exciting G333.6–0.2 contains stars that are very young, embedded in dust (Blum et al. 2002), in keeping with the youthful appearance of this extremely dense, compact H II region. On the other hand, the cluster exciting W43 must be older and more evolved because it contains a Wolf-Rayet star in addition to O giants or supergiants (Blum, Daminieli, & Conti 1999; Cotera & Simpson 1997), although there are still numerous stars forming in the associated molecular clouds (Motte, Schilke, & Lis 2003). This is also consistent with the more mature appearance of W43, where the ionized gas is well separated from the exciting star cluster.

In this study, we report measurements of the [N II] lines in the three positions observed in W43 by SCREH and numerous lines in three additional positions in G333.6–0.2 north of the center position observed by Colgan et al. (1993) with the KAO; we also report measurements of the central position in G333.6–0.2 with ISO’s Long Wavelength Spectrometer (LWS). In Section 2 we describe the observations, in Section 3 we discuss analysis techniques that are used to determine abundances in both low- and high-density H II regions, in Section 4 we model the two H II regions to determine the ICFs needed for those elements without measured abundances of the singly ionized atoms, and in Section 5 we discuss models in general and give our conclusions. By measuring both  $N^+$  and  $N^{++}$ , we hope to avoid the need for large (and uncertain) ICFs for nitrogen and thus improve the reliability of the secondary/primary abundance ratio measurement in the Milky Way molecular-ring H II regions.

## 2. Observations

Both G333.6–0.2 and W43 are extended compared to our 45–60" KAO beam and 80" LWS beam, which means that integrating over a constant density model to calculate ICFs is not adequate. Instead, we measured the fluxes at a number of positions at varying distances from the exciting stars to compare with the radial distributions of line flux computed from models to determine the source excitation and abundances.

### 2.1. KAO Observations

The observations were made with the Cryogenic Grating Spectrometer (CGS, Erickson et al. 1985) on NASA’s 0.91 m-KAO telescope. Data acquisition and reduction techniques are described by Colgan et al. (1993) and SCREH. For the present observations, Saturn was used for calibrating the absolute fluxes and determining the relative detector responses. The observed line and continuum fluxes for all positions are given in Table 1.

G333.6–0.2 was measured at its center and in three offset positions. The data from the center position, C (RA  $16^{\text{h}}22^{\text{m}}9^{\text{s}}.3$ , Dec  $-50^{\circ}6'5''$ , both J2000), were taken with a 45" beam and were previously reported by Colgan et al. (1993). The three North positions (N1, N2, and N3 at 1', 2', and 3' north of G333.6–0.2C) were measured on 1993 May 11, 13, and 20, flying from Christchurch, New Zealand. For these positions, a 60" beam was used for all lines to provide good ratios to the [N II] 205  $\mu\text{m}$  line, at whose wavelength the KAO has a diffraction size of  $1.22\lambda/D = 57''$ . The chopper angle was approximately east-west (position angle  $\sim 90^{\circ}$ ) and the throw was  $\sim 8.5'$ . Because the [N II] 122  $\mu\text{m}$  line is on the edge of a strong telluric  $\text{H}_2\text{O}$  line, the observations of positions N1, N2, and the [N II] 205  $\mu\text{m}$  line in Position N3 were taken at or near 43,000 feet altitude, instead of the more usual 39,000–41,000 feet, where the other lines in Position N3 were measured. Corrections for telluric  $\text{H}_2\text{O}$  absorption were made for both source and calibrator — we estimate the zenith  $\text{H}_2\text{O}$  overburden to range from 4.4 to 5.2  $\mu\text{m}$  at 43,000 feet on May 11 and 13, and 3.9  $\mu\text{m}$  at 41,000 feet on May 20.

The three positions of W43 that had previously been studied by SCREH (W43N at RA  $18^{\text{h}}47^{\text{m}}36^{\text{s}}.8$ , Dec  $-1^{\circ}55'29''$ ; W43C at RA  $18^{\text{h}}47^{\text{m}}36^{\text{s}}.0$ , Dec  $-1^{\circ}56'39''$ ; W43S at RA  $18^{\text{h}}47^{\text{m}}38^{\text{s}}.5$ , Dec  $-1^{\circ}57'34''$ , all J2000) in the [N III], [O III], and [S III] lines were measured in the [N II] lines on 1992 July 1, flying from NASA Ames. The 122  $\mu\text{m}$  line was measured in both 45" and 60" beams and the 205  $\mu\text{m}$  line was measured in a 60" beam. The chopper angle was approximately east-west, the throw was  $\sim 6'$ , and the zenith telluric  $\text{H}_2\text{O}$  overburden was  $\sim 5 \mu\text{m}$  at 43,000 feet and  $\sim 2 \mu\text{m}$  at 45,000 feet. We estimate that there

are extra measurement uncertainties of order 20% in addition to the statistical errors given in Table 1 for the W43 [N II] 122  $\mu\text{m}$  lines. These could result from clumpiness, pointing differences, or the effects of poor correction for deep atmospheric absorption. The analyses of the various measurements produce different abundance measures, as will be tabulated in the next sections.

## 2.2. ISO Observations

G333.6–0.2 was measured with the ISO Long Wavelength Spectrometer in LWS01 mode on 1997 March 20 (TDT No. 49001002, PI R. Rubin), exposure time 1054 s. The line strengths for the [O III] 52 and 88  $\mu\text{m}$  and [N III] 57  $\mu\text{m}$  lines are given in Table 1 (the [O I] 63  $\mu\text{m}$  line is also strong, but the [N II] 122  $\mu\text{m}$  line was not detected owing to ISO’s low spectral resolution in LWS01 mode with the LW detectors).

## 2.3. Extinction Corrections

The measured line fluxes were corrected for extinction using the infrared extinction function of Li & Draine (2001). This function has a very similar wavelength dependence to the extinction function used by SCREH when normalized to Li & Draine’s extinction maximum at  $\sim 9.5 \mu\text{m}$ . Thus the FIR optical depths are characterized as being proportional to the maximum optical depth in the 9.5 to 10  $\mu\text{m}$  region,  $\tau_{9.7}$ . Based on new information in the literature, we have updated these values from SCREH, with W43 having more extinction and G333.6–0.2 having less.

For three stars in the central cluster of W43, Blum et al. (1999) estimated that  $A_K = 3.55$ , 3.52, and 3.63. Using the extinction law of Li & Draine (2001), we find that  $A_K = 3.55$  corresponds to  $\tau_{9.7} = 2.74$ . Thus we increase the values of  $\tau_{9.7}$  used by SCREH by 37% to  $\tau_{9.7} = 2.74$ , 2.74, and 3.63 for W43N, W43C, and W43S, respectively.

For G333.6–0.2, Fujiyoshi et al. (2001) inferred an average value of  $\tau_{9.7} = 1.5$  from their mid-IR spectropolarization measurements. Consequently, we use  $\tau_{9.7} = 1.5$  for the center position instead of SCREH’s  $\tau_{9.7} = 2.0$ . (This value corresponds to  $A_{H\alpha} = 14$  mag; a lower extinction yet for the center of G333.6–0.2 might be predicted, since the nebula was detected in  $H\alpha$  by Churms et al. 1974.) We use  $\tau_{9.7} = 1.5$  for the North positions as well, since there are no significant variations in the integrated H I optical depths measured by Forster et al. (1987) at the C, N1, and N2 positions.

Following SCREH, an extinction uncertainty estimate of 20% ( $1\sigma$ ) is propagated through

all subsequent data analysis errors.

### 3. Analytical Approach

Because recombination, bremsstrahlung, and collisional excitation rates are all proportional to the electron density squared,  $N_e^2$ , it has long been the practice to treat all abundance analyses in H II regions and planetary nebulae (PNe) as though the surface brightness is, in fact, proportional to the emission measure,  $EM = \int N_e^2 dl$ , where  $l$  is the line-of-sight path through the nebula. This is valid so long as  $N_e$  is substantially less than the critical density,  $N_{crit}$ , defined as the density at which the upper level of a transition is depopulated equally by radiative transitions and by collisions (e.g., Osterbrock 1989; Rubin 1989). Corrections for collisional de-excitation can be made if  $N_e$  is known and approximately constant. However, constant densities are unlikely to occur in real nebulae. The effect of density variations is that the kernels in the integral over the nebular volume have different weighting functions for ions with different amounts of collisional de-excitation. Thus the flux ratios of such ions depend on these variable weighting functions (see, e.g., Rubin 1989). Here we consider the calculation of ratios of column densities of ions whose emission is proportional to  $N_e$ , as well as to  $N_e^2$ . The former is appropriate when  $N_e \gg N_{crit}$ . These two methods give identical results for ion ratios when the density is constant over the whole observed volume. We will demonstrate the effects of density variations via a simple example and use the results to estimate how to handle the transition case where  $N_e \sim N_{crit}$ .

#### 3.1. Low Density H II Regions

For a given ion, a “low density” H II region is one where  $N_e \ll N_{crit}$  and a “high density” H II region is one where  $N_e \gg N_{crit}$ ; thus any particular H II region can have both low density and high density volumes, depending on the value of  $N_{crit}$  for the transition being studied. Rubin (1989) tabulated values of  $N_{crit}$  for a large number of optical and IR transitions and Martín-Hernández et al. (2002a) tabulated values for many IR transitions; these values, of course, need to be updated as the atomic data are revised. The values of  $N_{crit}$  have a strong dependence on wavelength (owing to the dependence on wavelength of the transition probabilities), ranging from  $\sim 40 \text{ cm}^{-3}$  for the [N II] 205  $\mu\text{m}$  line to  $> 10^8 \text{ cm}^{-3}$  for some ultraviolet transitions. Because of their high  $N_{crit}$  compared to typical densities observed in H II regions and PNe, most optical and near-infrared (NIR) forbidden line and recombination line transitions are analyzed in the low density limit.

In the optically thin case (no dependence on line profile function), the line emissivity,  $j$ , used in the radiative transfer equation can be written as

$$j = \frac{h\nu}{4\pi} A \frac{N_{upper}}{N_i} N_i, \quad (1)$$

where  $A$  is the Einstein transition probability,  $N_i$  is the density of atoms in the particular ionization state,  $N_{upper}/N_i$  is the fraction of ions in the upper level of the transition, and  $h\nu$  is the energy of the transition. Then the observed line flux (ignoring extinction and self-absorption) is  $F = \int j d\Omega$  where  $d\Omega$  is the volume element of the H II region in the telescope beam,  $\Omega$ . We define the normalized volume emissivity  $\epsilon = 4\pi j/N_i N_e$  (Rubin et al. 1994; SCREH) and the average  $N_e$ -weighted ionic density as  $\langle N_i \rangle = \int N_i N_e d\Omega / \int N_e d\Omega$ . For the low density regime where  $N_e \ll N_{crit}$ ,  $\epsilon$  has very little dependence on  $N_e$ , and the line surface brightnesses (or fluxes) are proportional to  $\int N_i^2 dl$  (or  $\int N_e^2 d\Omega$ ).

Thus for this case we can write

$$F = \int \frac{\epsilon}{4\pi} N_i N_e d\Omega \approx \frac{\epsilon}{4\pi} \langle N_i \rangle \int N_e d\Omega \quad (2)$$

and

$$\frac{F_1}{F_2} = \frac{\epsilon_1}{\epsilon_2} \frac{\langle N_{i_1} \rangle}{\langle N_{i_2} \rangle}, \quad (3)$$

so the ionic abundance ratio is

$$\frac{\langle N_{i_1} \rangle}{\langle N_{i_2} \rangle} = \frac{F_1/\epsilon_1}{F_2/\epsilon_2}. \quad (4)$$

As  $N_e$  approaches  $N_{crit}$ , collisional de-excitation of the line's upper level becomes more and more important, and the ratio of 2 lines from the same ionic species  $N_i$  with different  $N_{crit}$  can be used to estimate the electron density (e.g., Rubin 1989; Fig. 1a).

### 3.2. High Density H II Regions

At high densities, the fractional population of the upper level approaches that of local thermodynamic equilibrium (LTE) and is a function of electron temperature,  $T_e$ , only. In this regime, the line surface brightness is proportional to the column density,  $CD = \int N_i dl$ , the number of ions in the line of sight:

$$F = \int \frac{(N_e \epsilon)}{4\pi} N_i d\Omega \approx \frac{(N_e \epsilon)}{4\pi} \Omega \int N_i dl, \quad (5)$$

where the quantity  $N_e \epsilon$  is the emissivity per ion and has very little density dependence so long as there are no low density regions along the line of sight (no large density variations).



Then for identical beam areas  $\Omega$ , the flux ratio is

$$\frac{F_1}{F_2} = \frac{(N_{e1}\epsilon_1)CD_1}{(N_{e2}\epsilon_2)CD_2}, \quad (6)$$

and the column density ratio is

$$\frac{CD_1}{CD_2} = \frac{F_1/(N_{e1}\epsilon_1)}{F_2/(N_{e2}\epsilon_2)}. \quad (7)$$

In fact, for abundance ratio calculations, one really wants to measure the actual number of ions in the line of sight, and *not* the number of ions weighted by some other, strongly varying parameter such as the electron density (as in equation (4)). Thus it is desirable to estimate column densities for gaseous nebulae whenever feasible.

### 3.3. Density Variations: An Illustrative Model

The presence of density variations has a number of effects on the abundances that one infers from FIR line measurements. We demonstrate these effects by the use of simple slab models, with an exponentially decreasing density  $N_e = N_0 e^{-x}$  from a location in the interior of the slab  $x_0$  to an outer location  $x_{max}$ . This density function was chosen because it gives similar weighting to both high and low density regimes (other types of models, e.g., Rubin 1989, could be used to illustrate the same effects). For simplicity, the example assumes no ionization stratification,  $N_i = N_e$ , and  $T_e = 7000$  K; thus the only effects on the calculated volume emissivities are those of the variation in density. For all models, a very large value of  $x_{max}$  was chosen such that  $N_e = 0.01 \text{ cm}^{-3}$ . Each model had a different value of  $x_0$  and  $N_0$ , ranging from  $N_0 = 1 \text{ cm}^{-3}$  to  $N_0 = 10^6 \text{ cm}^{-3}$ . The intensities emitted in each line were calculated from

$$I = \int_{x_0}^{x_{max}} \frac{\epsilon(x)}{4\pi} N_i N_e dx. \quad (8)$$

Figure 1a shows how the electron density,  $N_{est}$ , can be estimated using calculated volume emissivity ratios from the observed [S III] 19/33  $\mu\text{m}$ , [O III] 52/88  $\mu\text{m}$ , and [N II] 122/205  $\mu\text{m}$  line ratios, with the [N II] ratio being sensitive to low densities and the [S III] ratio sensitive only to much higher densities. Values of  $N_{est}$  were computed from the ratios of the line intensities produced by the models; ratios of these estimated densities to the maximum density of each model,  $N_0$ , are plotted in Figure 1b. This figure shows that in the presence of density variations, even when the ions coincide exactly in space, one does not infer the same density from the three ratios:  $N_{est}(\text{N II}) < N_{est}(\text{O III}) < N_{est}(\text{S III})$ . From this, one must conclude that the volumes containing the observed ions are not co-extensive (as is the

case for, e.g., ionization stratification) if the estimated densities are in the range of density sensitivity and do not follow these inequalities (subject to the observational and atomic data uncertainties).

To understand the effects of density variations on estimated abundances, we calculate the estimated numbers of ions in the volume from the model line intensities and  $N_{est}$  and compare these estimates with the true numbers of ions found by integrating  $N_e$  or  $N_e^2$  over the model line of sight. These comparisons we call "abundance measures". Numerically, we define the abundance measures,  $AM$ , as either the calculated line intensity  $I$  divided by the normalized volume emissivity  $\epsilon$  for low density analysis or  $I$  divided by the emissivity per ion ( $N_{est}\epsilon$ ) for high density analysis. These quantities are normalized by dividing by the appropriate integrals of density squared or density, respectively. That is, the abundance measure for analysis that assumes that the line intensities are proportional to the emission measure is

$$AM_{EM} = \frac{I}{\epsilon} \left[ \int_{x_0}^{x_{max}} N_i N_e dx \right]^{-1} \quad (9)$$

and the abundance measure for analysis that assumes that the line intensities are proportional to the column density is

$$AM_{CD} = \frac{I}{(N_{est}\epsilon)} \left[ \int_{x_0}^{x_{max}} N_i dx \right]^{-1}. \quad (10)$$

$AM_{EM}$  and  $AM_{CD}$  are plotted in Figure 2 for each ion as functions of  $N_{est}$ . Both abundance measures contain fewer ions than the input number, but the regimes of reasonable validity for the two measures are different and depend on  $N_{crit}$ . Fortunately, much of the discrepancy with the  $N_e$  or  $N_e^2$  volume integrals cancels when we take ratios of the  $EM$  or  $CD$  abundance measures (equivalent to equations (4) and (7), respectively), as shown in Figure 3. These plots suggest an approach for the intermediate cases where  $N_{est} \sim N_{crit}$ .

First, there must be measurements of  $N_e$  from (at least) two different density-sensitive line pairs:  $N_{e1}$  and  $N_{e2}$ . Then, if the column density of the more density-sensitive ion of one pair is the numerator and that of the less density-sensitive ion from the other pair is the denominator, the column density ratio is greater than the true value and the  $N_e^2$ -weighted ionic ratio is less than the true value. This follows directly from equations (2)-(7), from which we derive

$$\frac{CD_1}{CD_2} = \frac{< N_{i1} > N_{e2}}{< N_{i2} > N_{e1}}. \quad (11)$$

Likewise, if the less density sensitive ion of one pair is the numerator and the more density sensitive ion from the other pair is the denominator, the column density ratio is less than the true value and the  $N_e^2$ -weighted ionic ratio is greater than the true value.

We note that because  $N_{crit}$  for the N III 57  $\mu\text{m}$  line is intermediate between the values of  $N_{crit}$  for the O III 52 and 88  $\mu\text{m}$  lines, the estimated  $\text{N}^{++}/\text{O}^{++}$  ratio has little uncertainty for all values of  $N_e$ .

Looking at Fig. 3, we define the three regimes as follows: The low-density regime occurs when the estimated density is lower than the lowest  $N_{crit}$  for the lines involved (see Fig. 2 for  $N_{crit}$ ) — here we use the equations for emissivity  $\propto N_e^2$  since the ionic ratios weighted by  $N_e^2$  are close to the model input ratio (unity for the plots). The high-density regime occurs where the column density ratio is closer to the model input ratio than the ionic ratios weighted by  $N_e^2$  — this occurs when the estimated density is higher than the larger  $N_{crit}$  of the more density-sensitive line pair. At intermediate densities, both the low-density and high-density equations give abundances discrepant from the model input. Considering this, in the intermediate regime we average the two cases and quote an uncertainty equal to the greater of one half the difference between the two abundance ratios and the statistical error. (If one knew for sure that the volumes containing the different ions were co-extensive, then an elaborate weighting scheme could be devised, but the overall uncertainties do not warrant such schemes.) For the low- and high-density regimes, we will use the equations for the ionic ratios weighted by  $N_e^2$  or the column density ratios, respectively. Unfortunately, it is extremely difficult to get a good estimate of the column density of  $\text{H}^+$  because the emissivities of H lines in H II regions and PNe are always proportional to  $N_e^2$  and the properly weighted density is impossible to determine. Because of this, we are not able to measure abundance ratios with respect to hydrogen for all the ions; instead we will obtain standard “ionic ratios” with respect to  $\text{H}^+$  only for those ions whose lines have high  $N_{crit}$  compared to the densities estimated from the observed line pairs. We will calculate both ionic ratios and column density ratios for the rest of the heavy element ions observed at FIR wavelengths and consider the density regime according to the line’s  $N_{crit}$  for estimating the final abundance ratio.

A significant difference between this method and previous abundance determinations from FIR lines is the increase in the size of the errors that occur when more than one density-sensitive line pair is measured in a high-density, compact H II region. *The smaller statistical errors that one obtains when only one density line pair is measured are an illusion.* It is clear that measuring multiple density-sensitive line pairs is not just desirable but is, in fact, required to obtain reliable abundance measurements from FIR lines (even though they may have large assigned errors). We note that a similar situation occurs in optical abundance determinations from forbidden lines in that  $T_e$  must be inferred from various line ratios to obtain reliable abundances (e.g., Kennicutt, Bresolin, & Garnett 2003).

## 4. G333.6-0.2 and W43

### 4.1. Analytical Abundance Estimates

The complete set of derived  $N_e$  values and abundance ratios is given in Table 2. The values of  $T_e$  used were 6200 K (Caswell & Haynes 1987) and 6500 K (Wink, Wilson, & Bieging 1982) for G333.6-0.2 and W43, respectively (mid-and FIR lines are not sensitive to uncertainties in  $T_e$ ). The radio fluxes for G333.6-0.2 were estimated by integrating over the model described in the next section at the appropriate distances from the center (Figure 4) and the radio fluxes for W43 were taken from Lester et al. (1985) (the NRAO VLA Sky Survey, Condon et al. 1998, has lower fluxes, probably because the bright clumps in W43 are somewhat optically thick at 1.4 GHz, and the “snapshot” observation of Balser et al. 1995 is missing extended flux). The differences between the ionic abundance results presented here and in SCREH for the KAO data are due to different extinction corrections (§2.3), different radio fluxes for G333.6-0.2, and different collisional excitation cross sections for [S III] (Tayal & Gupta 1999). The differences between the ISO and KAO abundances with respect to  $H^+$  probably indicate the uncertainties in the radio flux estimates. The  $N^{++}/O^{++}$  ratios are the same to within the errors ( $3\sigma$ ), although some of the difference may be due to different beam sizes or pointing errors.

Estimated abundance ratios for N/S and S/H are given in Table 3, where we assumed that the ionization fraction for sulfur  $\langle S^{++}/S \rangle = 0.9$  (a typical value for the low-to-moderate excitation and high metallicity models in Rubin 1985, SCREH, and the models computed in §5) and that the ionization fraction of  $\langle N^{3+}/N \rangle$  is so small that  $N^{3+}$  can be neglected. The N/H ratio was derived by multiplying N/S times S/H. Abundance estimates for the other elements are more uncertain because they require the use of models to estimate the necessary ICFs.

### 4.2. Models of G333.6-0.2 and W43

We computed H II region models using the program NEBULA (Rubin 1968a; Rubin 1985; SCREH). We have upgraded NEBULA since the discussion in SCREH in the following ways:

(1) The photoionization and recombination cross sections were extensively updated to include results from the Opacity Project (Seaton et al. 1992). Fits to photoionization cross sections by Verner et al. (1996) were used for photoionization of  $C^+$ ,  $C^{++}$ ,  $N^0$ ,  $N^+$ ,  $N^{++}$ ,  $O^0$ ,  $O^+$ ,  $Ne^0$ ,  $Ne^+$ ,  $S^+$ ,  $S^{++}$ ,  $S^{3+}$ , and  $Si^+$ . The changes to the coefficients for ionization

and recombination of iron are discussed by Rodríguez & Rubin (2004). (Changes to other elements were omitted because of the lack of corresponding Opacity Project recombination coefficients.) Total recombination coefficients (including dielectronic recombination) were taken from the tabulations of Nahar & Pradhan (1997) for  $C^+$ ,  $C^{++}$ , and  $N^{++}$ ; Kisielius & Storey (2002) for  $N^+$ ; Nahar (1999) for  $O^0$  and  $O^+$ ; Kisielius et al. (1998) for  $Ne^+$ ; Nahar (2000) for  $S^{++}$ ; and Nahar (1995) for  $S^+$  and  $Si^+$ . Charge exchange rate coefficients for reactions with  $H^+$  were taken from the tabulation of Kingdon & Ferland (1996) with updates from the Oak Ridge National Laboratory web site “Charge Transfer Database”<sup>3</sup> and coefficients for reactions with  $He^+$  were taken from J. G. Wang et al. (2003, in preparation), with the data available from the same Oak Ridge web site.

(2) The collisional excitation cross sections now include  $N^+$  and  $O^{++}$  from Lennon & Burke (1999),  $N^{++}$  from Blum & Pradhan (1992),  $Ne^+$  from Saraph & Tully (1994),  $Ne^{++}$  from McLaughlin & Bell (2000),  $S^{++}$  from Tayal & Gupta (1999),  $S^{3+}$  from Saraph & Storey (1999) or Tayal (2000) for  $T_e > 10,000$  K,  $Ar^+$  from Pelen & Berrington (1995), and  $Ar^{++}$  from Galavís, Mendoza, & Zeippen (1995, 1998).

(3) The diffuse radiation field calculation (Rubin 1968a) was modified to include calculations at all radial grid positions instead of interpolation between selected positions. Up to 20 iterations were used if needed to obtain uniform convergence (most positions converged within 10 iterations).

(4) Non-LTE stellar atmospheres with winds from Pauldrach, Hoffmann, & Lennon (2001) were used. Smith et al. (2002) have also computed model atmospheres using Pauldrach’s code. They note that models computed with these atmospheres give much better results for computations of starburst galaxy spectra than the atmospheres from Schaerer & de Koter (1997), which were used in the H II region models of Stasińska & Schaerer (1997) (we did not use Smith et al.’s models because of the coarseness of their wavelength grid). Lumsden et al. (2003) also show that the Pauldrach et al. (2001) atmospheres give better results for comparison of measurements of helium lines in Galactic H II regions.

We used the set of model atmospheres, which were downloaded from Pauldrach’s web site<sup>4</sup>, consisting of models for Solar composition, dwarfs (D) and supergiants (S), and  $T_{\text{eff}} = 30, 35, 40, \text{ and } 45 \times 10^3$  K. Although G333.6–0.2 and W43 have metallicities somewhat larger than that of the Orion Nebula, the Orion Nebula abundances (e.g., Simpson et al. 1998) are sufficiently below the “Solar” composition used in the stellar atmosphere models

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<sup>3</sup><http://www-cfadc.phy.ornl.gov/astro/ps/data/>

<sup>4</sup><http://www.usm.uni-muenchen.de/people/adi/adi.html>

that G333.6–0.2 and W43 are probably closer in metallicity to the “Solar” composition models than to anything greatly different (recent work finds a lower Solar O/H ratio, as will be discussed later). Stellar atmospheres with higher metallicities have more winds, which produce more extreme ultraviolet (EUV) flux, but also more line blanketing (less EUV); the difference is probably mainly for frequencies  $> 4$  Ryd, which are not important for H II regions. We computed models for W43 using both D and S models from Pauldrach et al. (2001) and also the non-LTE, plane-parallel model atmospheres without winds calculated by Lanz & Hubeny (2003)<sup>5</sup>. We used only the dwarf star model atmospheres for G333.6–0.2 because it is clearly excited by a very young star cluster (Blum et al. 2001),

One result of the changes to the nebular atomic physics data is that the volume-averaged ionization fractions (SCREH; Rubin et al. 1994) for  $\langle \text{O}^{++}/\text{O} \rangle$  and  $\langle \text{Ne}^{++}/\text{Ne} \rangle$  predicted by the models are larger than previous predictions, whereas the fractions for  $\langle \text{S}^{++}/\text{S} \rangle$  and  $\langle \text{N}^{++}/\text{N} \rangle$  are approximately the same. The differences are similar to typical uncertainties in ionization ratios computed from observational data ( $\lesssim 30\%$ ), which would make testing difficult unless the data sample was large and of very high quality. However, the choice of stellar atmosphere model makes a much larger difference in the ionization fractions and ratios thereof; this choice has serious implications for estimates of  $T_{\text{eff}}$  and ICFs. We will compare various volume-averaged ionization fractions computed for different model H II regions as a function of stellar atmosphere model and discuss the need for winds further in §5.1.

#### 4.2.1. G333.6–0.2

The initial model is that of a spherically symmetric, core-envelope H II region (Rubin, Hollenbach, & Erickson 1993), with density decreasing from high values in the core to low values in the outer envelope, as suggested by the estimated densities of Table 2. Our initial density distribution for the G333.6–0.2 model computed here was derived by inverting the radio brightness temperature distribution, since the radio surface brightness is proportional to  $\int N_e^2 dl$ . The radio brightness temperature distribution of G333.6–0.2, of necessity, has been estimated by combining a variety of different measurements, none of which includes the whole H II region at a frequency high enough to insure that the emission is everywhere optically thin. The highest spatial resolution radio maps in the literature consist of the following maps: (1) a Parkes 64-m telescope map at 8.87 GHz with half-power beamwidth, HPBW =  $2.5'$ , made by McGee, Newton, & Butler (1979); because the map was made with a single dish, it covers the full spatial extent of G333.6–0.2 to a radius of  $> 15'$ ; (2) a Fleurs

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<sup>5</sup> Available from <http://tlsty.gsfc.nasa.gov>

Synthesis Telescope map at 1.415 GHz with HPBW  $\sim 50''$  by Retallack & Goss (1980); only the inner  $5'$  could be mapped, and the core is optically thick; (3) an Australia Telescope Compact Array map at 3.4 cm (8.8 GHz) with resolution  $1.6''$  by Fujiyoshi (1999); only the inner  $13''$  could be mapped, and the peak brightness temperature of  $\sim 4400$  K indicates that the central  $2'' - 3''$  is not completely optically thin ( $T_e \sim 6200$  K, Caswell & Haynes 1987). However, this last map is extremely important because it shows just how compact and bright the central core of G333.6-0.2 really is. A compact bright core is also required by the NIR maps and the density-sensitive FIR line ratios.

Still lacking is a measurement of the transition region between the inner core and the envelope. There is, however, a map of the warm dust emission made by the Midcourse Space Experiment (MSX; Price et al. 2001) at  $21 \mu\text{m}$  with  $19''$  resolution.<sup>6</sup> Although there is no certainty that warm dust has the same spatial structure as ionized gas (and it certainly does not have constant temperature), in the absence of an alternative we used it to interpolate between the core and envelope radio maps.

The total radio flux was taken from the measurements at the highest observed frequencies because G333.6-0.2 shows evidence of being at least partially optically thick at frequencies  $< 10$  GHz. We used the integrated flux of 93.8 Jy at 14.7 GHz measured by McGee & Newton (1981). The final estimated radio brightness distribution was renormalized to produce this total integrated radio flux. The estimated radio brightness temperature distribution, scaled  $\propto \nu^{-2.1}$  as though it were optically thin at 5 GHz, and derived electron densities are shown in Figure 4. The total radio flux requires a stellar ionizing photon luminosity of  $10^{50}$  photons  $\text{s}^{-1}$ . There may also be, of course, additional EUV photons that do not contribute to the H II region photoionization because they are absorbed by dust. Table 4 lists the relevant IR and radio observations, corrected for extinction, including those mid-IR observations from the literature which include most of the core.

In addition to matching the radio brightness temperature distribution, the observations of G333.6-0.2 put a number of requirements on the parameters of any model. There must be a core of very high density to produce the bright radio and mid-IR core and the high [O III]  $52/88 \mu\text{m}$  and [S III]  $19/33 \mu\text{m}$  line ratios seen in the center position. The material at  $2' - 4'$  must be of low density to produce the  $52/88 \mu\text{m}$  and  $19/33 \mu\text{m}$  line ratios of order unity ( $N_e$  of order  $150 - 400 \text{ cm}^{-3}$ ). The material further out must be of even lower density ( $N_e \sim 40 \text{ cm}^{-3}$ ) to make a large contribution to the  $122/205 \mu\text{m}$  line ratio  $\sim 2.5$ . The stellar  $T_{\text{eff}}$  has to be high enough to produce  $\text{O}^{++}$  and  $\text{N}^{++}$  at all positions (out to  $> 3'$ ) but low enough so that there is ample  $\text{N}^+$ ,  $\text{Ar}^+$ , and not much  $\text{N}^{++}$ ,  $\text{O}^{++}$ , or  $\text{S}^{3+}$  even in the

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<sup>6</sup>Available at <http://irsa.ipac.caltech.edu>

center position. In addition, the  $\text{He}^+/\text{H}^+$  radio recombination line (RRL) ratio is only 0.045 (McGee & Newton 1981), indicative of a relatively low  $T_{\text{eff}}$  ( $\lesssim 35,000$  K).

We computed over a hundred models of G333.6–0.2 with a variety of combinations of stellar atmosphere models and density filling factors, all models reproducing the observed radio flux. The abundance ratios used in the models were estimated from the fluxes in Table 4 (see also Table 3) and range from 1.25 to 3.4 times the Orion Nebula abundances (Simpson et al. 1998). Line fluxes were computed from the models by integrating over flat-topped circular beam profiles for the IR lines or circular Gaussian beam profiles for the radio recombination lines located at the projected distances of the observed positions from the H II region center.

Our best example of a spherically-symmetric model is the G333.6–0.2 Core model in Table 4. This particular model is shown because it predicts line ratios that agree reasonably well with the density indicators ( $[\text{S III}]$  19/33  $\mu\text{m}$ ,  $[\text{O III}]$  52/88  $\mu\text{m}$ , and  $[\text{N II}]$  122/205  $\mu\text{m}$ ) and the ionization indicator measured in the largest beam ( $\text{He}^+/\text{H}^+$  RRLs measured by McGee & Newton 1981). The model uses a linear combination of 0.9375 times the D-30 (30,000 K Dwarf) plus 0.0625 times the D-35 stellar atmosphere models of Pauldrach et al. (2001) and has a volume filling factor of 0.5. The contribution from the D-35 model, although seemingly small, is necessary to have enough high energy photons to ionize He and N to the observed ionization. The given model combination only represents the shape of the ionizing spectrum and not the actual numbers or spectral types of the ionizing stars.

However, no model that we computed is a good fit — the data seem to require stellar ionizing fluxes with both higher and lower average  $T_{\text{eff}}$ . The large fluxes of low-excitation ions, such as  $\text{Ne}^+$  and  $\text{Ar}^+$  in the center of G333.6–0.2 seem to require stars with low  $T_{\text{eff}}$  ( $\lesssim 32,000$  K was suggested by Rank et al. 1978). However, the central density and optical depth are so high that when the exciting star is this cool, all the high energy photons that are capable of ionizing  $\text{O}^+$  to  $\text{O}^{++}$  are absorbed at small radii (0.87 pc, corresponding to 60") and the models have essentially no  $\text{O}^{++}$  at the radii corresponding to our three North positions. Yet we observe  $\text{O}^{++}$  at all four positions. It would seem that an exciting star with a somewhat higher  $T_{\text{eff}}$  ( $> 35,000$  K) is required or a combination of stars including a much hotter star; however, for such models the  $> 35$  eV photons required to ionize  $\text{O}^+$  to  $\text{O}^{++}$  still are confined in the central core but the total  $[\text{O III}]$  line fluxes are much larger than observed unless the O/H, O/N, O/Ne, etc. ratios are unreasonably low. On the other hand, in H II regions ionized by stars with relatively low  $T_{\text{eff}}$ , the  $\text{N}^{++}$  Strömgren radius,  $R_S$ , is much larger than the  $\text{O}^{++}$   $R_S$ , leading to observed  $\text{N}^{++}/\text{O}^{++}$  ratios that can be much larger than the actual N/O ratio (e.g., Rubin et al. 1994; SCREH; Stasińska & Schaerer 1997). The result of the differences in the  $\text{N}^{++}$  and  $\text{O}^{++}$   $R_S$  is that the observed  $\text{N}^{++}/\text{O}^{++}$  ratio should



increase with projected distance from the center. This predicted increase is not observed in our  $N^{++}$  and  $O^{++}$  data taken at the North positions. Likewise, the ISO data in a  $80''$  beam taken at the center position do not produce a larger  $N^{++}/O^{++}$  ratio than does our KAO data in a  $45''$  beam. Again, the observed behavior of this line ratio can be reproduced only by models with higher  $T_{\text{eff}}$  stars. In fact, all the models that we computed predict a much larger  $[N \text{ III}] 57/[O \text{ III}] 52 \mu\text{m}$  ratio than is observed at any position. Agreement cannot be improved by increasing the O/H ratio (the models already predict more  $[O \text{ III}] 52 \mu\text{m}$  flux than is observed) or decreasing the N/H ratio (the model  $[N \text{ II}]$  lines fluxes agree with the observations). A lower  $T_{\text{eff}}$  seems to be required, but already there is too little  $O^{++}$  in the North positions. We conclude that a spherical model, even of varying density, cannot satisfy the observations of both the very bright core and the doubly ionized lines in the low density region at  $2' - 4'$  from the center.

A possible solution might be some sort of non-spherical model with large clumps, or a blister model (face-on from its symmetrical appearance, Hyland et al. 1980). The NEBULA code for a blister model (Rubin et al. 1991) assumes that the density distribution is that of the surface of a molecular cloud with some possible gradient of density going into the cloud. The stellar photons ionize a hole on the edge of the cloud. However, that axisymmetric model does not use a density function that allows for a particularly high density immediately next to the stars, as is required for a model of G333.6-0.2. Without a modification to the axisymmetric 2-D code, this code is not suitable. A treatment that can handle a two- or three-dimensional density distribution appears to be required.

We next considered a quasi-blister model or two-component model, such as was used by Simpson et al. (1986) and Colgan et al. (1991) to model the Orion Nebula and K3-50, respectively. This would represent either a blister H II region with one open side or an H II region with very large clumps where ionizing photons escape through holes or paths between the clumps. A linear combination was made of two spherical models, one with the high density core and one with an almost empty center to provide the doubly ionized ions seen at  $2' - 4'$ . The model with the empty center had a very large  $R_S$  and therefore had to be arbitrarily truncated in order that it be density bounded at about the same radius as the  $R_S$  for the Core model. An example of such a shell model, labeled "Model for the North Positions", is also given in Table 4; the model used the same stellar atmosphere as the Core model but the filling factor was unity. Even here there are not enough photons to ionize  $O^{++}$  at the N2 and N3 positions and too much  $N^{++}$  at the N1 position. Possibly three or even four spherical models with different interior shell diameters and different scaling fractions for the volume of the model in the line of sight are needed to reproduce the observations at the four observed positions. Because of the arbitrariness of the procedure and the non-uniqueness of the result, we did not calculate any of these multi-shell models.

The abundances have been only partially adjusted to produce the observed fluxes — the immediate goal of the models was to determine  $T_{\text{eff}}$ , the filling factor, and other geometry factors. Further adjustment to the abundances is not meaningful because the geometry still is not well determined, although the agreement of the  $\text{He}^+/\text{H}^+$  RRL ratios with the observations indicates that  $T_{\text{eff}}$  is not especially discrepant (since most helium is primordial, the interstellar  $\text{He}/\text{H}$  ratio is not increased significantly as a result of Galactic chemical evolution). The uncertainty in the geometry factors makes it difficult to estimate a reliable abundance for oxygen. However, if we consider just the Center and N1 positions, we find that the observed fluxes would be matched by the model fluxes if  $\text{O}/\text{H} \sim 3 \times 10^{-4}$ , a value somewhat lower than that observed in the Orion Nebula. This  $\text{O}/\text{H}$  abundance ratio is not reliable, but suggests that  $\text{O}/\text{H}$  is not especially high in G333.6–0.2. As expected, since we observed the dominant ionization states of nitrogen and sulfur ( $\text{N}^+$  and  $\text{S}^{++}$ ) in this source, our final estimated abundances of these elements are those input to the models (Table 3). On the other hand, it is difficult to think of a geometry that would produce the large  $[\text{Ar II}]$   $6.98 \mu\text{m}$  line flux observed in the G333.6–0.2 core, given the number of high energy photons that are needed to produce the observed  $[\text{Ne III}]$  and  $[\text{O III}]$  lines. This discrepancy may indicate that the atomic data for argon need revision (see also Morisset et al. 2003).

Kim & Koo (2001) showed that many, if not most, ultracompact H II regions are not ionization bounded at the radius given by the ultracompact component, but have extended envelopes ionized by the same stars. They suggested that the ultracompact components are blister H II regions and the extended envelopes are the result of champagne flows combined with hierarchical clumping. Although G333.6–0.2 is larger than the typical ultracompact H II region because of its multiple stars and high ionizing luminosity, it does have the high central density and optical depth of an ultracompact H II region. We conclude that G333.6–0.2 may be an example of a large ultracompact H II region with an extended envelope.

#### 4.2.2. W43

Because the constant density model by SCREH for W43 is not a good match to the densities inferred from the observations (see Table 2), we calculated new models for it. The appearance of W43 is that of a large “T” or incomplete “D” shape projected on the sky (Lester et al. 1985; Liszt 1995; Two Micron All Sky Survey image<sup>7</sup>) with a central star cluster near but not on the vertical element of the letter. Because our three observed

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<sup>7</sup>2MASS, <http://irsa.ipac.caltech.edu>

positions, including the one closest to the cluster, all have similar densities and excitation (SCREH; Table 2), we interpret the geometry as a cluster of ionizing stars surrounded by a cavity and then a shell of ionized gas, such that all three positions are at essentially the same physical distance from the stars. Thus we model the H II region as a shell of ionized gas with interior radius = 1.2 pc (0.75' at a distance of 5.6 kpc (Reifenstein et al. 1970; Liszt 1995), scaled to a Galactic Center distance of 8 kpc). The model  $N_e$  drops smoothly from  $770 \text{ cm}^{-3}$  at 1.2 pc to  $420 \text{ cm}^{-3}$  at  $R_S \sim 2.9$  pc, matching the average of the total observed flux of 77.5 Jy at 2 cm (Schraml & Mezger 1969) and 97.4 Jy at 6 cm (Reifenstein et al. 1970) for an ionizing luminosity of  $2.9 \times 10^{50} \text{ photons s}^{-1}$ . The densities and ionic abundance ratios of the three positions were averaged (weighted) for comparison with the models; these comparisons are given in Table 5. Three models were computed with the only difference being the input stellar atmosphere models, as described in the Table. The reason the S (supergiant) model atmospheres were also used is the presence in W43 of a Wolf-Rayet star and two O giants/supergiants (Blum et al. 1999). The combinations of stellar atmosphere models were adjusted until the H II region models produced the observed average radio  $\text{He}^+/\text{H}^+$  ratio ( $\text{He}/\text{H} = 0.10$  for the nebular models) and the observed  $\text{N}^{++}/\text{N}^+$  ratio.

To correct the observed abundance ratios for unseen ionization states, we used the ionization fractions calculated from the model as shown in the Table 5 comments and footnotes. The biggest difference between the models is for the  $\langle \text{O}^{++}/\text{O} \rangle / \langle \text{S}^{++}/\text{S} \rangle$  and  $\langle \text{N}^{++}/\text{N} \rangle / \langle \text{O}^{++}/\text{O} \rangle$  ratios, no doubt because the Supergiant model atmospheres have less flux from 35 to 41 eV than the Dwarf model atmospheres. Pauldrach et al.'s (2001) Supergiant models have much higher mass loss rates than their Dwarf models, so the difference must be caused by more line blanketing at these energies by the gas in the wind. On the other hand, the static, plane-parallel, non-LTE atmospheres of Lanz & Hubeny (2003), which have no winds, have much lower fluxes from  $\sim 40$  to  $\sim 50$  eV.

All three W43 models adequately reproduce the  $T_{\text{eff}}$ -indicating indices: the  $\text{He}^+/\text{H}^+$  and  $\text{N}^{++}/\text{N}^+$  ratios. The range of implied abundances in Table 5, though, indicates the large uncertainties resulting from determination of abundances through model-produced ICFs, particularly when different stellar atmosphere models are used to compute the ICFs. We prefer the values of  $\text{N}/\text{H}$  derived from comparisons with sulfur, rather than those derived from oxygen. The estimated  $\text{N}/\text{O}$  ratio is 0.32 to 0.42 from ICFs computed with models D and S (not very different from the 0.34 estimated by SCREH), but is as low as 0.16, similar to that of the Orion Nebula (Rubin et al. 1991), from the ICFs computed with model T. The reason for the wide range is not the abundances of nitrogen or sulfur, which are reliable because they do not require large ICFs. The problem is the total abundance of oxygen: the model ICFs produce  $\text{O}/\text{H}$  ranging from  $2.8$  to  $8.3 \times 10^{-4}$ , whereas the predicted  $\text{O}/\text{H}$  ratio is  $7.2 \times 10^{-4}$  for a Galactic abundance gradient of  $d(\log \text{O}/\text{H})/dR = -0.06 \text{ dex kpc}^{-1}$  (Henry

& Worthy 1999), given the Orion Nebula’s location at  $R_G \sim 8.4$  kpc and  $O/H = 4.0$  to  $4.4 \times 10^{-4}$  (Rubin et al. 1991; Esteban et al. 1998). In fact, one expects that the oxygen abundance for W43 should be much higher than that of the Orion Nebula, given W43’s low  $T_e = 5410$  to  $6500$  K (Subrahmanyan & Goss 1996; Wink, Wilson, & Bieging 1983; and references in both papers). For this source, as for G333.6–0.2, we conclude that we still do not have a good estimate for the abundance of oxygen, the most abundant heavy element and the most important contributor to cooling in H II regions.

## 5. Discussion and Conclusions

### 5.1. Stellar Atmosphere Models

As was discussed above, many of the ionization correction factors are critically dependent on the choice of stellar atmosphere model. This can be demonstrated either by plotting ratios of the numbers of photons able to ionize various ions as a function of  $T_{\text{eff}}$  or by plotting ionization fractions or ICFs calculated from models; here we do the latter because ICFs are more readily applicable to abundance computation. This is illustrated in Figure 5, where we plot the ICF ratios  $\langle N^{++}/N \rangle / \langle O^{++}/O \rangle$ ,  $\langle O^{++}/O \rangle / \langle S^{++}/S \rangle$ , and  $O^{++}/O$  versus the abundance-independent excitation ratio  $\langle N^{++}/N \rangle / \langle N^+/N \rangle$ . The models from Rubin (1985) that are plotted in Fig. 5 all have total nucleon density  $= 10^3 \text{ cm}^{-3}$ , ionizing luminosities of  $10^{49}$  to  $10^{50}$  photons  $\text{s}^{-1}$ , and Orion Nebula metallicities; the models from SCREH had similar densities and ionizing luminosities but metallicities ranging from 1 to 2 times Orion nebula metallicities. The models from both Rubin (1985) and SCREH were computed with the versions of NEBULA available at the time and both used LTE stellar atmosphere models from Kurucz (1979). In addition to the models described in §4, constant nucleon density ( $= 10^3 \text{ cm}^{-3}$ ) spherical models with ionizing luminosities of  $10^{49}$  photons  $\text{s}^{-1}$  were computed and plotted. These recent NEBULA models used stellar atmosphere models from either Pauldrach et al. (2001), Lanz & Hubeny (2003), or contemporary<sup>8</sup> LTE stellar atmosphere models described by Kurucz (1992). The stellar atmosphere models had “Solar” abundances and the nebular models had 1.5 times Orion Nebula abundances (Simpson et al. 1998) such that  $O/H = 6.0 \times 10^{-4}$ . Because they are widely used, we also plot ICFs from the models of Stasińska & Schaerer (1997). For these models both the stellar and nebular abundances were “Solar”, with  $O/H = 8.51 \times 10^{-4}$  (currently it is believed that the Solar  $O/H$  ratio is lower:  $5.45 \times 10^{-4}$ ; Holweger 2001).

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<sup>8</sup><http://kurucz.harvard.edu/>

Unless an H II region can be observed at both optical and infrared wavelengths, observers cannot measure the abundances of N, O, Ne, S, and Ar without making corrections for unseen ionization states for some elements. For optically obscured H II regions like W43 and G333.6–0.2, the most significant missing ion is  $O^+$ , since oxygen is the most abundant heavy element and since these low excitation H II regions are believed to have most of their oxygen as  $O^+$ . Rubin et al. (1994) and SCREH assumed that they knew the O/S ratio and used the observed  $O^{++}/S^{++}$  ratio to estimate the ICF for their measured  $N^{++}/O^{++}$  ratios, since they did not have any other line pair that could be used to indicate excitation. In this paper we use our observed  $N^{++}/N^+$  ratios to estimate  $T_{\text{eff}}$ ; however, we see from Fig. 5a how dependent the  $N^{++}/O^{++}$  ICF is on the stellar atmosphere model.

Martín-Hernández et al. (2002a) took advantage of their ISO SWS measurements of the  $Ne^{++}/Ne^+$  ratios to estimate ICFs without the need for assuming *a priori* some abundance ratio; unfortunately, the large beam size difference between the SWS and LWS instruments means that their observed  $Ne^{++}/Ne^+$  ratios may not always be appropriate for estimating the excitation of  $O^{++}$  and  $N^{++}$ . Moreover,  $Ne^{++}$  requires ionizing photons  $> 41$  eV, and it is at these high energies that the stellar atmosphere models are especially divergent, as we now discuss.

An important test of the validity of stellar atmosphere models of hot stars is whether H II region models produced with these atmospheres predict line fluxes that agree with observations. Except for a few rare instances, this is the only way that the EUV fluxes of the models can be tested, since interstellar absorption by hydrogen predominantly prevents the EUV fluxes from hot O stars from reaching the Earth. Here we discuss the stellar fluxes that can doubly ionize neon ( $h\nu > 41$  eV). This has been called the [Ne III] problem (e.g., Sellmaier et al. 1997) because it has been observed that the  $Ne^{++}/O^{++}$  ratio is relatively constant over a large range of H II region excitation (SCREH; Stasińska & Schaerer 1997 and references therein; Kennicutt et al. 2003), contrary to predictions of models that used LTE stellar atmospheres (e.g., SCREH). Proposed solutions have involved using non-LTE stellar atmospheres, usually with winds (Rubin, Kunze, & Yamamoto 1995; Sellmaier et al. 1996; Stasińska & Schaerer 1997). Figure 6 shows the observations of SCREH along with the predictions of the nebular models discussed in this paper. Models calculated with black bodies for the stellar atmosphere spectrum are also plotted since Morisset et al. (2003) found that such models give a reasonable fit to the mid-IR ISO observations.

Surprisingly, in Fig. 6 only the models of Stasińska & Schaerer (1997) reproduce the observations over the total range of observed  $O^{++}/S^{++}$ , although if Ne/O is as large as 0.25 in all H II regions as it is in the Orion Nebula (Simpson et al. 1998; Table 3), these H II region models produce too high a  $Ne^{++}/O^{++}$  ratio because of the very large EUV fluxes in

Schaerer & de Koter’s (1997) stellar atmosphere models. Moreover, Schaerer & de Koter’s models have too much EUV flux for reproducing other observations (e.g., Smith et al. 2002). Earlier, Sellmaier et al. (1997) had demonstrated that the [Ne III] problem was supposedly solved when they obtained a good fit to the data by using non-LTE atmospheres with winds computed with Pauldrach’s code as it then existed. However, our H II region models with non-LTE stellar atmospheres with winds from Pauldrach et al. (2001) taken from both A. Pauldrach’s web site and from the stellar atmosphere models produced by Smith et al. (2002) using Pauldrach’s WM-basic code predict much lower  $\text{Ne}^{++}/\text{O}^{++}$  than observed for Dwarf atmospheres with  $T_{\text{eff}} < 40,000$  K and for Supergiant atmospheres with  $T_{\text{eff}} < 35,000$  K. Our H II region model produced with a 35,000 K supergiant atmosphere from F. Sellmaier (private communication in 1995) lies above the Pauldrach et al. Supergiant line but not as high as the model with the 35,000 K supergiant atmosphere in Sellmaier et al.’s (1996) Figure 2, even though we are now assuming  $\text{Ne}/\text{O} = 0.25$  (Simpson et al. 1998) instead of 0.2025. Part of the difference of the current nebular model predictions with previous models is the higher  $\text{O}^{++}$  ionization relative to  $\text{Ne}^{++}$  and  $\text{S}^{++}$  obtained with the Opacity Project cross sections described in Section 4. The model computed with Sellmaier’s (1995) atmosphere has higher ionization than the model computed with Pauldrach et al.’s (2001) S-35 atmosphere because it used Orion Nebula abundances (as in Sellmaier et al. 1996) instead of the higher metallicity (approximately 2 times the Orion Nebula abundances of Table 3) used for the rest of the models plotted here. The result, though, is a return of the mismatch between the stellar atmosphere models and the EUV fluxes of real stars, with possible additional discrepancies for the ionization, recombination, and collisional excitation cross sections for oxygen, neon, and sulfur. Moreover, we see that all the models computed without winds do a poorer job of reproducing the observations than both the models computed with non-LTE atmospheres with winds and the models computed with simple black bodies for stellar atmospheres (for this reason we give low weight to the T model for W43 in Table 5). This probably means that the codes producing the stellar atmosphere models used for computation of H II region models should include the effects of stellar winds.

## 5.2. Abundances and Abundance Gradients

At the low excitations seen in inner Galaxy H II regions, the predicted flux of the  $\text{O}^{++}$  lines is extremely sensitive to the stellar atmosphere, but if the models are accurate, the O/H ratio of both G333.6–0.2 and W43 could be as low as  $3 - 5 \times 10^{-4}$ , substantially lower than would be predicted from the observed Orion Nebula abundance and an O/H abundance gradient of  $-0.06$  dex  $\text{kpc}^{-1}$  (Henry & Worthey 1999). This low O/H ratio is hard to understand because of the low  $T_e$  measured in inner Galaxy H II regions from

radio recombination lines. There would need to be some source of systematic error for the O/H ratio to be higher than these estimates: i.e., what would be required is that the excitation of oxygen in the two H II regions is lower than estimated in this paper from  $\text{He}^+/\text{H}^+$  recombination lines and  $\text{N}^+$  and  $\text{N}^{++}$  forbidden lines. Possible reasons that the G333.6–0.2 and W43 excitation might be overestimated are as follows: (1) The abundance of  $\text{N}^+$  is underestimated because the [N II] line fluxes are undermeasured owing to telluric absorption at  $122\ \mu\text{m}$  or diffraction at  $205\ \mu\text{m}$ . This is a particular problem for W43, where the two measurements of the [N II]  $122\ \mu\text{m}$  line in different apertures can be used to produce very different electron densities and  $\text{N}^{++}/\text{N}^+$  abundance ratios (1.5 to 3.0 in Table 5). (2) The estimated  $\text{N}^{++}/\text{N}^+$  ratio is too high because some of the estimated  $\text{N}^+$  is missing owing to inadequate correction for density variations. (3) The estimated  $\text{N}^{++}/\text{N}^+$  ratio is too high because of incorrect collisional excitation cross sections for the pertinent energy levels. (4) The models predict too much  $\text{O}^{++}/\text{O}$  versus  $\text{N}^{++}/\text{N}^+$ . The spread of  $\text{O}^{++}/\text{O}$  ratios in Fig. 5c shows that any excessive  $\text{O}^{++}/\text{O}$  is due to model atmospheres and not due to the  $\sim 30\%$  higher oxygen ionization resulting from the Opacity Project cross sections. (5) The ionization of helium might be overestimated if the He/H abundance ratio is actually greater than the assumed 0.10. We conclude that determinations of the O/H abundance ratio in low excitation H II regions have large uncertainties at this time, with the largest contributor to the uncertainty being the choice of stellar atmosphere models, with some possible contribution from the atomic data used in the nebular models.

We have demonstrated that the abundances of nitrogen and sulfur can be obtained from FIR observations and the ratio of N/S can be obtained when both ionization states of  $\text{N}^+$  and  $\text{N}^{++}$  are measured. (Simpson et al. 1998, Martín-Hernández et al. 2002a, and Giveon et al. 2002 also show that the abundances of Ne/H and Ar/H, as well as S/H, can be determined from mid-IR observations with little bias owing to excitation or abundance.) Compared to SCREH, the S/H abundance ratio is higher in G333.6–0.2, owing to the new data at the additional North positions, but lower in W43 because of the higher collisional excitation cross sections for sulfur. The N/H ratio is about 50% larger in G333.6–0.2, owing to the new data at the North positions, but a factor of 2 smaller in W43, owing to the smaller ICF for  $\text{N}^+$ . However, since the estimated O/H ratio is also much smaller in W43, now that we are estimating the ionization from  $\text{N}^{++}/\text{N}^+$  and not from  $\text{O}^{++}/\text{S}^{++}$  with the assumption of a constant O/S ratio as SCREH did, the estimated N/O ratios agree with SCREH for both G333.6–0.2 and W43. Thus we conclude that with this new analysis, the S/H and N/O ratios are consistent with the abundance gradients measured by SCREH, Rudolph et al. (1997), and Afflerbach et al. (1997), and Martín-Hernández et al. (2002a) for N/O. The N/H ratio for G333.6–0.2 is also consistent with the gradient of  $-0.10\ \text{dex kpc}^{-1}$  estimated by SCREH and Afflerbach et al. (1997) but the N/H ratio inferred for W43 is lower than

what would be expected from the previously published abundance gradients.

On the other hand, the estimated O/H ratio is consistent with no gradient at all, working inward from the Orion Nebula. This is in contradiction to the other abundance studies of H II regions and PNe in the Galaxy, particularly in regions other than the inner Galaxy where optical studies allow measurement of  $O^+$  (see the reviews of Henry & Worthy 1999 and Rolleston et al. 2000 and references therein). Smartt et al. (2001) measured abundances in four B stars in the inner Galaxy ( $R_G \sim 2.5$  to 5 kpc) by optical absorption line spectroscopy. They found that these four stars have oxygen abundances similar to stars in the Solar neighborhood, even though they have enhanced abundances of nitrogen, magnesium, aluminum, silicon, and sulfur corresponding to their inner Galaxy locations and expected abundance gradients of  $-0.07$  dex  $kpc^{-1}$ . They have no explanation for this anomalous oxygen abundance, nor do we. We conclude that more measurements of the oxygen abundance in inner Galaxy stars and H II regions are needed to confirm this apparent lack of increase in the O/H ratio, because we know that oxygen is one of the main coolants in H II regions and the low observed radio  $T_e$  measurements require additional cooling from heavy elements compared to Solar-neighborhood H II regions.

### 5.3. Conclusions Regarding Systematic Errors

There are several possible sources of systematic error that must be considered when measuring abundances in obscured, low-excitation H II regions, that is, all the H II regions in the inner Galaxy. The impact of these systematic errors is that there is a very large uncertainty in the total metallicity and important abundance ratios like N/O. These sources of systematic error are the following:

(1) The most abundant heavy element ion in low-excitation H II regions,  $O^+$ , has no bright infrared lines. The consequence is that the abundance of this critically important element must be estimated from the abundance of an ionization state with fractional abundance  $< 0.5$ .

(2) Nitrogen, the secondary element which one needs for studies of galaxy chemical evolution, has two important ionization states:  $N^+$  as well as  $N^{++}$ . Fortunately, both species have measurable FIR lines, but *both* must be measured to determine the excitation.

(3) Moreover, both FIR  $N^+$  lines have greatly different values of  $N_{crit}$  from those of  $S^{++}$  or  $N^{++}$ . Thus one needs to be aware of and compensate for the ever-present density variations.



(4) The average density could be quite different in the various ionization zones, as shown by the different electron densities derived from  $N_e$ -diagnostic line pairs: low ionization ( $N^+$ ), intermediate ionization ( $S^{++}$ ), and high ionization ( $O^{++}$ ).

(5) ICFs estimated from models have major uncertainties owing to uncertainties in the stellar atmosphere models and possibly the atomic data used in the H II region models.

It is clear that the abundances of primary (oxygen, neon, or sulfur) and secondary (nitrogen) elements will need to be measured in many more inner Galaxy H II regions before we understand the chemical evolution of the Milky Way. In the near future, measurements of  $N^{++}$ ,  $N^+$ ,  $O^{++}$ ,  $Ne^{++}$ , and  $S^{++}$  will be possible from the Stratospheric Observatory for Infrared Astronomy, SOFIA. We recommend that future observers measure excitation-sensitive lines, such as  $Ne^+$  and  $Ne^{++}$ , as well as both  $N^+$  and  $N^{++}$  to determine the H II region ionization state needed to estimate the ICF for  $O^+$ . It is important that a sufficient number of lines be measured in each H II region (and extended H II regions mapped) so that detailed models of the H II region can be made. Discrepancies between the predicted line fluxes from the models and the observations can then be used to indicate EUV energy regions where the stellar atmosphere models may need revision.

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Fig. 1.— (a) Ratios of volume emissivities for line pairs of the electron-density-sensitive lines [S III] 19/33  $\mu\text{m}$ , [O III] 52/88  $\mu\text{m}$ , and [N II] 122/205  $\mu\text{m}$  are plotted as a function of electron density,  $N_e$ , for electron temperature,  $T_e$ , of 7000 K. (b) Line intensities were calculated for slab models with  $N_e = N_0 e^{-x}$  by integrating the volume emissivity of each line from  $x_0$  to  $x_{\text{max}}$ .  $N_0$  ranges from 1 to  $10^6 \text{ cm}^{-3}$  and  $N_e = 0.01 \text{ cm}^{-3}$  at  $x_{\text{max}}$  (see text). Values of  $N_e$  were estimated from the abscissa of Fig. 1a by assuming that the ratios of the line intensities could be used for the ordinate. These estimated values of  $N_e$  divided by  $N_0$  are plotted versus  $N_0$ . The root-mean-squared (RMS) density divided by  $N_0$  is also plotted (dotted line).

Fig. 2.— Line intensities were calculated for slab models with  $N_e = N_0 e^{-x}$  by integrating the volume emissivity of each line from  $x_0$  to  $x_{\text{max}}$ ; corresponding densities were estimated from the density-sensitive line ratios in Fig. 1a. The number of ions estimated from the line fluxes and estimated volume emissivities are plotted for lines of sight with varying amounts of high-density gas (see text); the abscissa is the estimated  $N_e$  for the line of sight. These abundance measures,  $AM_{EM}$  and  $AM_{CD}$ , are defined as the calculated line intensity divided by the normalized volume emissivity or the emissivity per ion, respectively. They are normalized by dividing by the emission measure or by the column density of ions, respectively (see text). Because preferred abundance measures for each ion depend on the density, critical densities,  $N_{\text{crit}}$ , are also tabulated for  $T_e = 7000 \text{ K}$ . (a)  $\text{S}^{++}$  measured from [S III] 19 and 33  $\mu\text{m}$  lines. (b)  $\text{O}^{++}$  and  $\text{N}^{++}$  measured from [O III] 52, [O III] 88, and [N III] 57  $\mu\text{m}$  lines. Note the near coincidence of the  $\text{N}^{++}$  and  $\text{O}^{++}$  abundance measures. (c)  $\text{N}^+$  measured from [N II] 122 and 205  $\mu\text{m}$  lines.

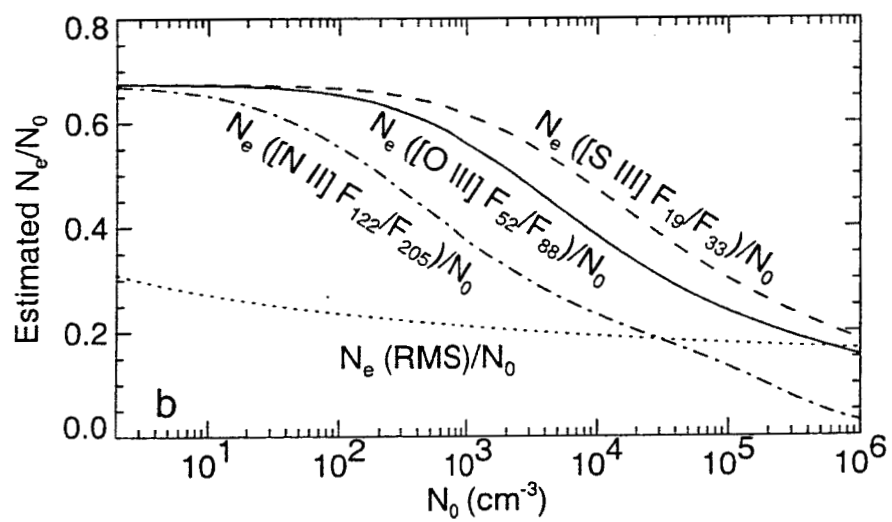
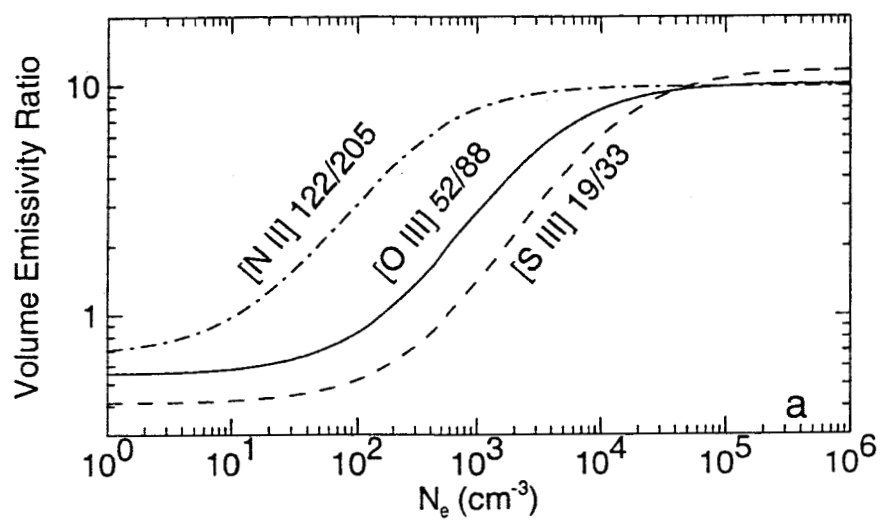
Fig. 3.— Ratios of the abundance measures for the ions shown in Fig. 2 are plotted as a function of estimated  $N_e$  for lines of sight with varying amounts of high-density gas. The heavy line is  $AM_{EM}$  and the dotted line is  $AM_{CD}$  (the light line is unity, for reference). (a)  $\text{N}^{++}/\text{O}^{++}$ . The lines for  $AM_{EM}$  and  $AM_{CD}$  coincide. (b)  $\text{O}^{++}/\text{S}^{++}$ . (c)  $\text{N}^+/\text{S}^{++}$ . (d)  $\text{N}^{++}/\text{S}^{++}$ . (e)  $\text{N}^{++}/\text{N}^+$ .

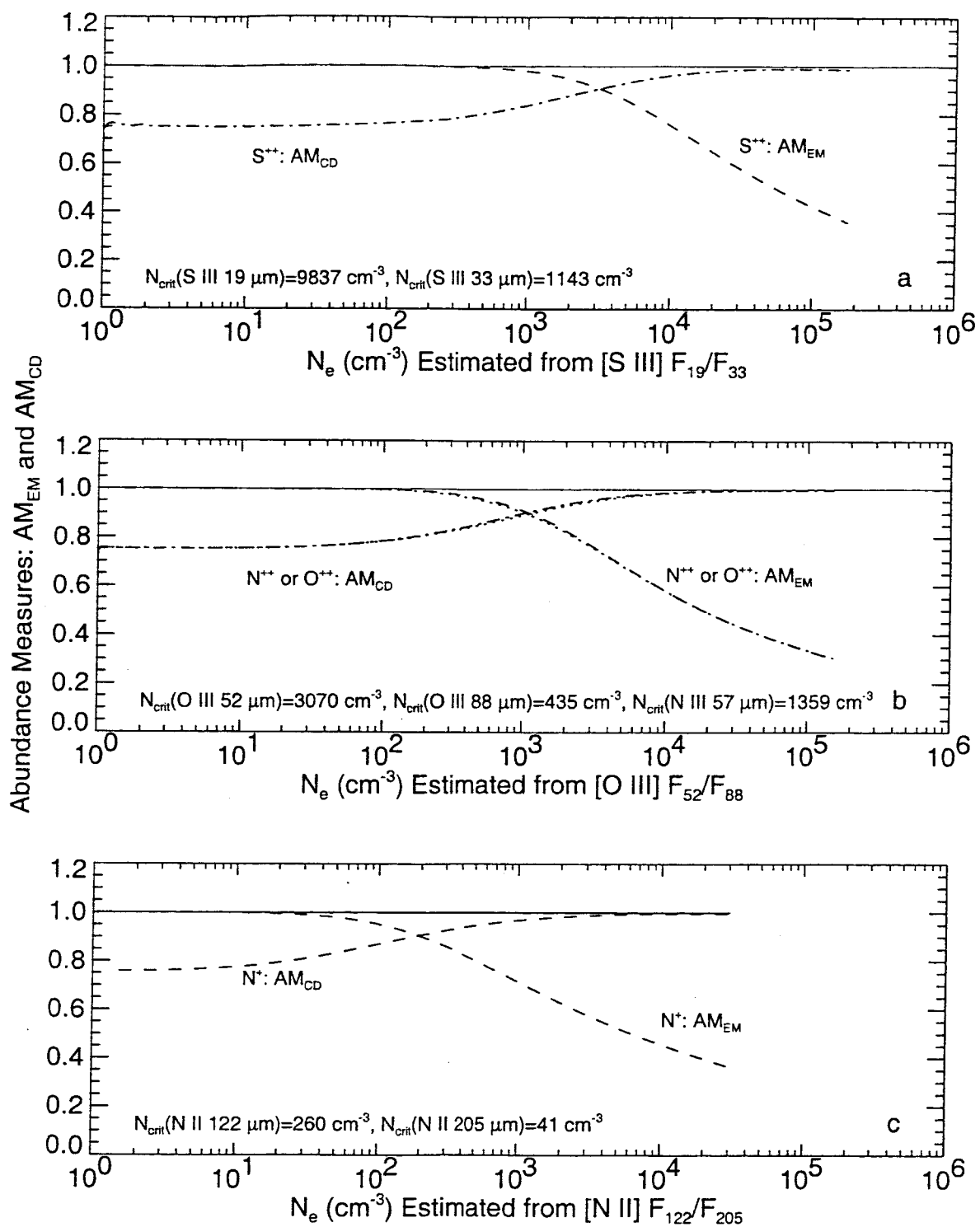
Fig. 4.— (a) Azimuthally-averaged radio brightness temperatures for G333.6-0.2 are plotted as a function of angular distance on the sky from the central peak. The radio brightness temperatures were all scaled to the same frequency, 5 GHz, using the formulation of Mezger & Henderson (1967) and ignoring optical depth effects. (b)  $N_e$ , derived by inverting the radio brightness temperature distribution shown in Fig. 4a, is plotted as a function of distance,  $r$ , from the central star cluster. A spherical H II region at a distance of 3 kpc was assumed. The overall form of this density function is  $N_e \propto r^{-1.2}$ , as shown by the dashed line, which was fit by least squares.

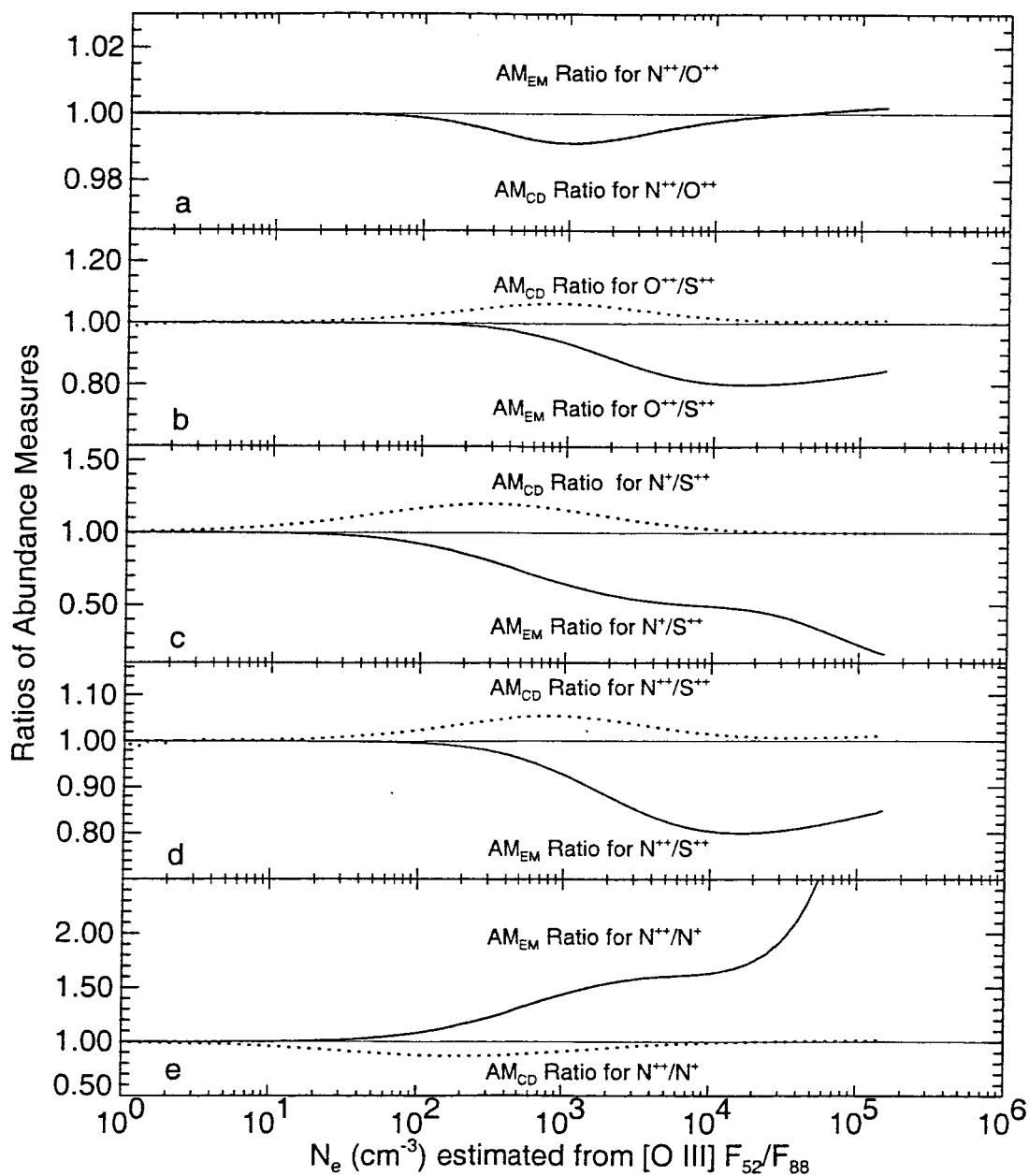


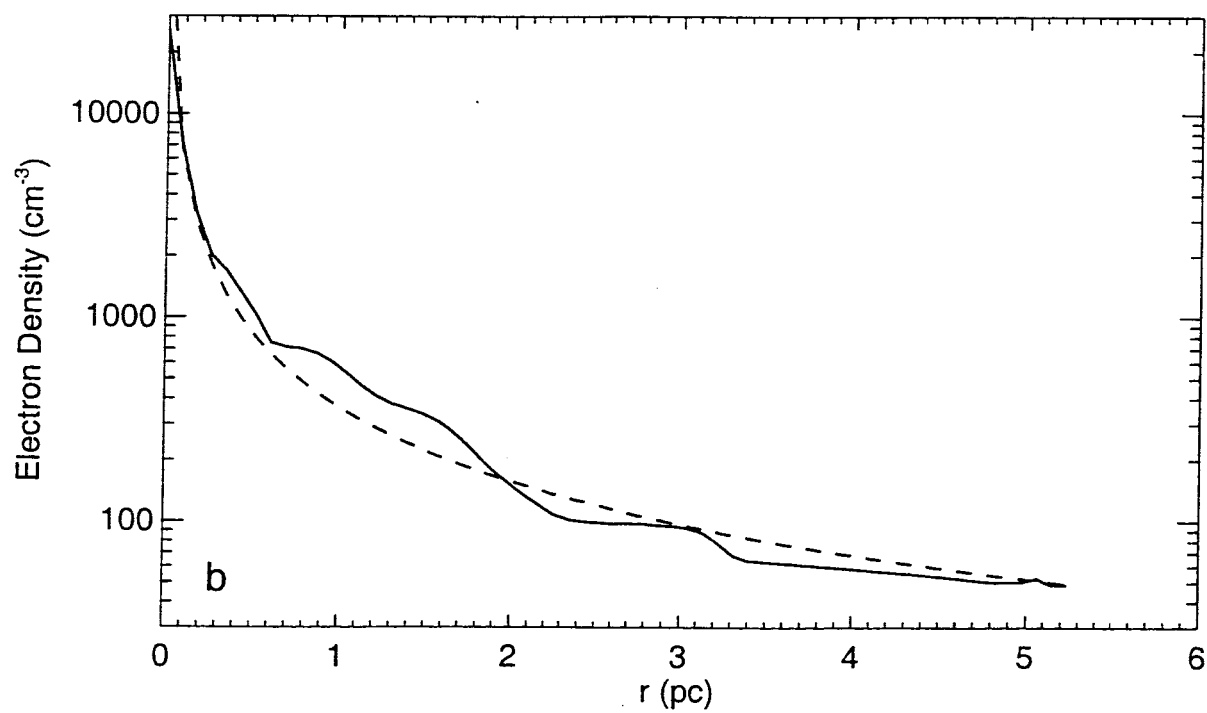
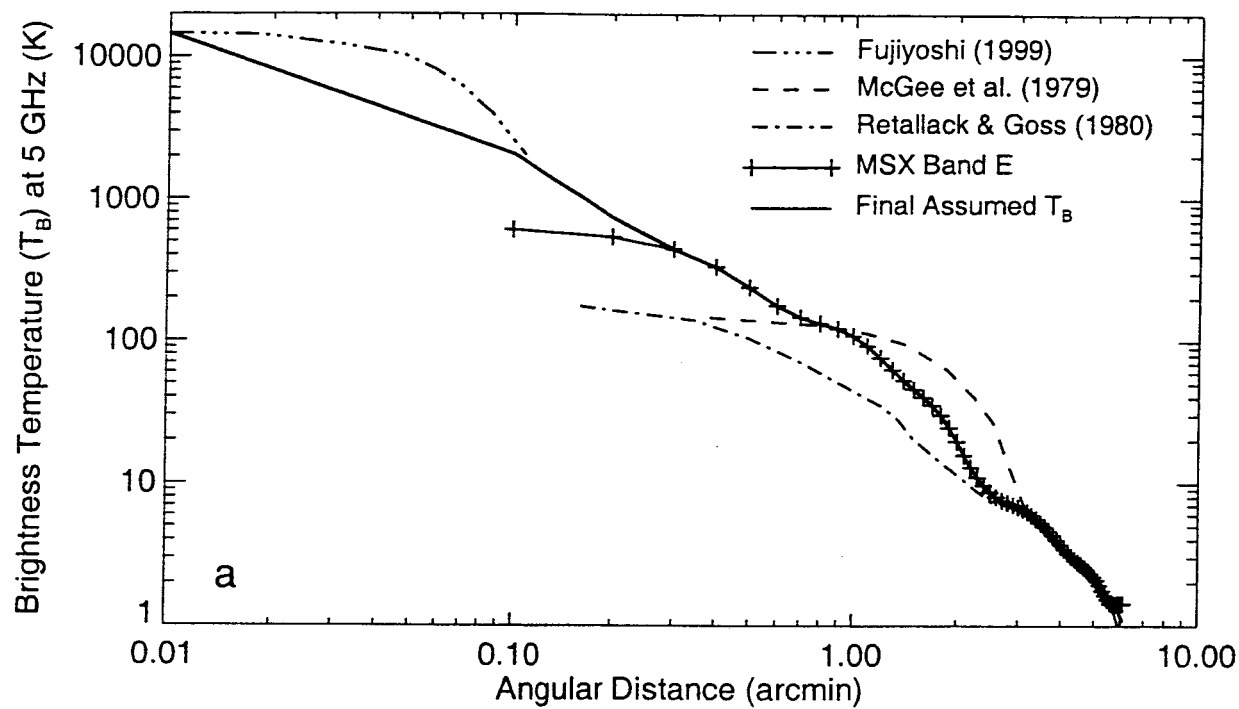
Fig. 5.— Ratios of ionization fractions are plotted for spherical H II region models. The models with LTE stellar atmosphere models are those from Rubin (1985) and SCREH, who used model atmospheres from Kurucz (1979), and the new NEBULA models with contemporary atmosphere models described by Kurucz (1992) ( $T_{\text{eff}} = 33, 35, 37, 40$ , and  $45 \times 10^3$  K). The H II region models with non-LTE stellar atmosphere models are those from Stasińska & Schaerer (1997), who used atmosphere models from Schaerer & de Koter (1997), and the NEBULA models calculated for this paper, where we used non-LTE atmosphere models from Pauldrach et al. (2001) and Lanz & Hubeny (2003) ( $T_{\text{eff}} = 30, 35, 40$ , and  $45 \times 10^3$  K). For the H II region models calculated with Pauldrach et al. atmospheres, the solid line connects models with Dwarf atmospheres and the dashed line connects models with Supergiant atmospheres. For the H II region models calculated with Lanz & Hubeny atmospheres, the solid line connects models with atmospheres with  $\log g = 4.0$  and the dotted and dashed lines connect models with atmospheres with  $\log g = 3.0$  to  $3.5$  and with Lyman continuum luminosities of  $10^{49}$  and  $10^{50}$  photons  $\text{s}^{-1}$ , respectively. The notation in the legend applies to all parts of the figure. (a)  $\langle \text{N}^{++}/\text{N} \rangle / \langle \text{O}^{++}/\text{O} \rangle$  versus  $\langle \text{N}^{++}/\text{N} \rangle / \langle \text{N}^+/\text{N} \rangle$ . (b)  $\langle \text{O}^{++}/\text{O} \rangle / \langle \text{S}^{++}/\text{S} \rangle$  versus  $\langle \text{N}^{++}/\text{N} \rangle / \langle \text{N}^+/\text{N} \rangle$ . (c)  $\langle \text{O}^{++}/\text{O} \rangle$  versus  $\langle \text{N}^{++}/\text{N} \rangle / \langle \text{N}^+/\text{N} \rangle$ .

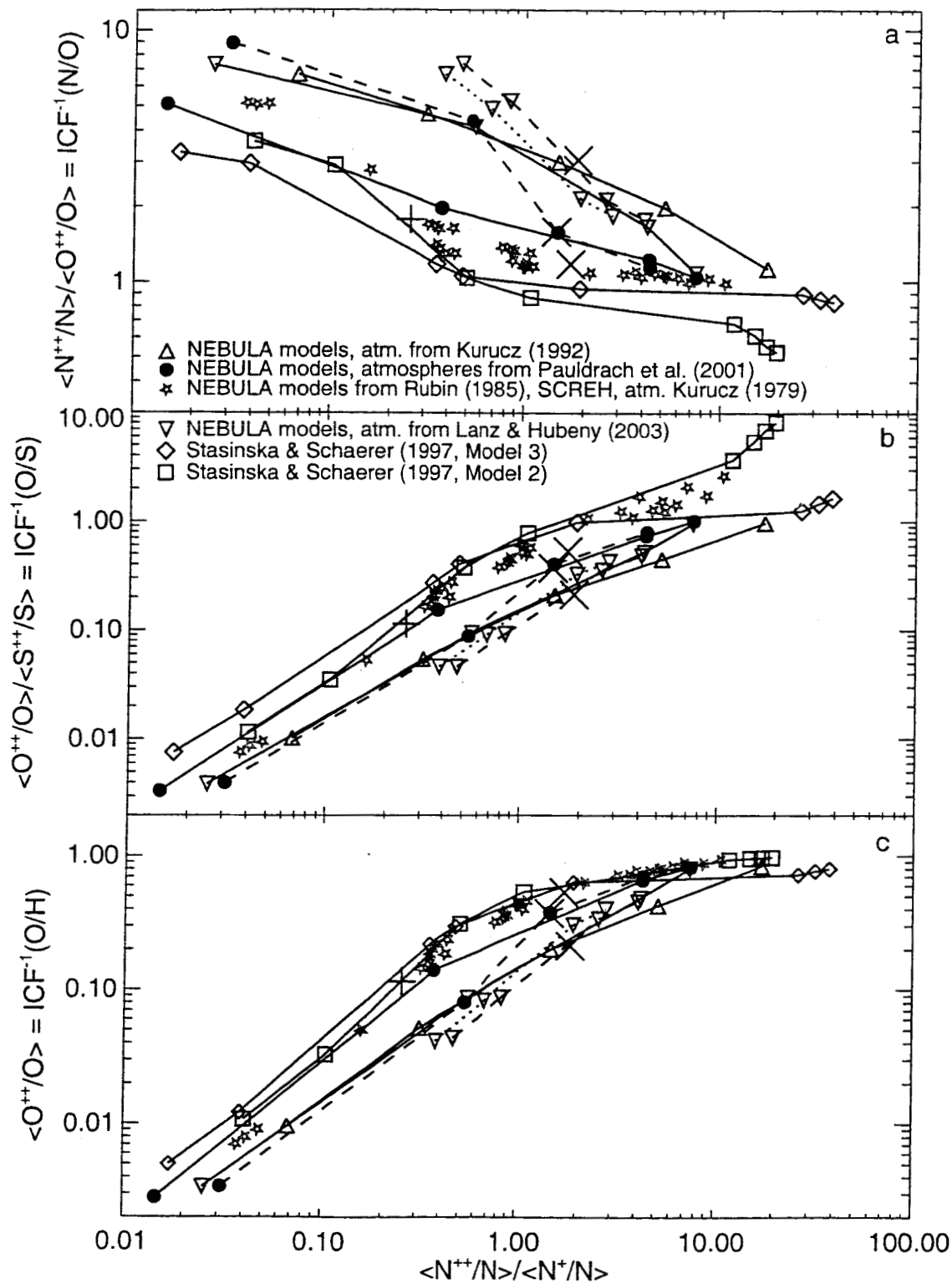
Fig. 6.— Ratios of  $\text{Ne}^{++}/\text{O}^{++}$  inferred from the observations of Simpson et al. (1995) are plotted against their inferred  $\text{O}^{++}/\text{S}^{++}$  ratios. Ratios from the nebular models discussed in this paper are also plotted. The notations of Pauldrach et al. (2001), Sellmaier (1995), Lanz & Hubeny (2003), and Kurucz (1992) in the legend refer to the stellar atmospheres used in our NEBULA models (same  $T_{\text{eff}}$  as the models in Fig. 5; Sellmaier’s atmosphere is a Supergiant with  $T_{\text{eff}} = 35,000$  K). The plotted  $\text{Ne}^{++}/\text{O}^{++}$  ratio for the models is equal to the calculated  $\langle \text{Ne}^{++}/\text{Ne} \rangle / \langle \text{O}^{++}/\text{O} \rangle$  ratio multiplied by an assumed abundance ratio of  $\text{Ne}/\text{O} = 0.25$  and the plotted  $\text{O}^{++}/\text{S}^{++}$  ratio for the models is equal to the calculated  $\langle \text{O}^{++}/\text{O} \rangle / \langle \text{S}^{++}/\text{S} \rangle$  ratio multiplied by an assumed abundance ratio of  $\text{O}/\text{S} = 50$  (Orion Nebula abundance ratios from Simpson et al. 1998). For the H II models calculated with Pauldrach et al. atmospheres, the solid line connects models with Dwarf atmospheres and the dashed line connects models with Supergiant atmospheres. For the H II region models calculated with Lanz & Hubeny atmospheres, the solid line connects models with atmospheres with  $\log g = 4.0$  and the dotted and dashed lines connect models with atmospheres with  $\log g = 3.0$  to  $3.5$  and with Lyman continuum luminosities of  $10^{49}$  and  $10^{50}$  photons  $\text{s}^{-1}$ , respectively. Models computed with atmospheres represented by black bodies from 20,000 K to 40,000 K (every 5000 K) are also plotted (open circles connected by a dotted line).











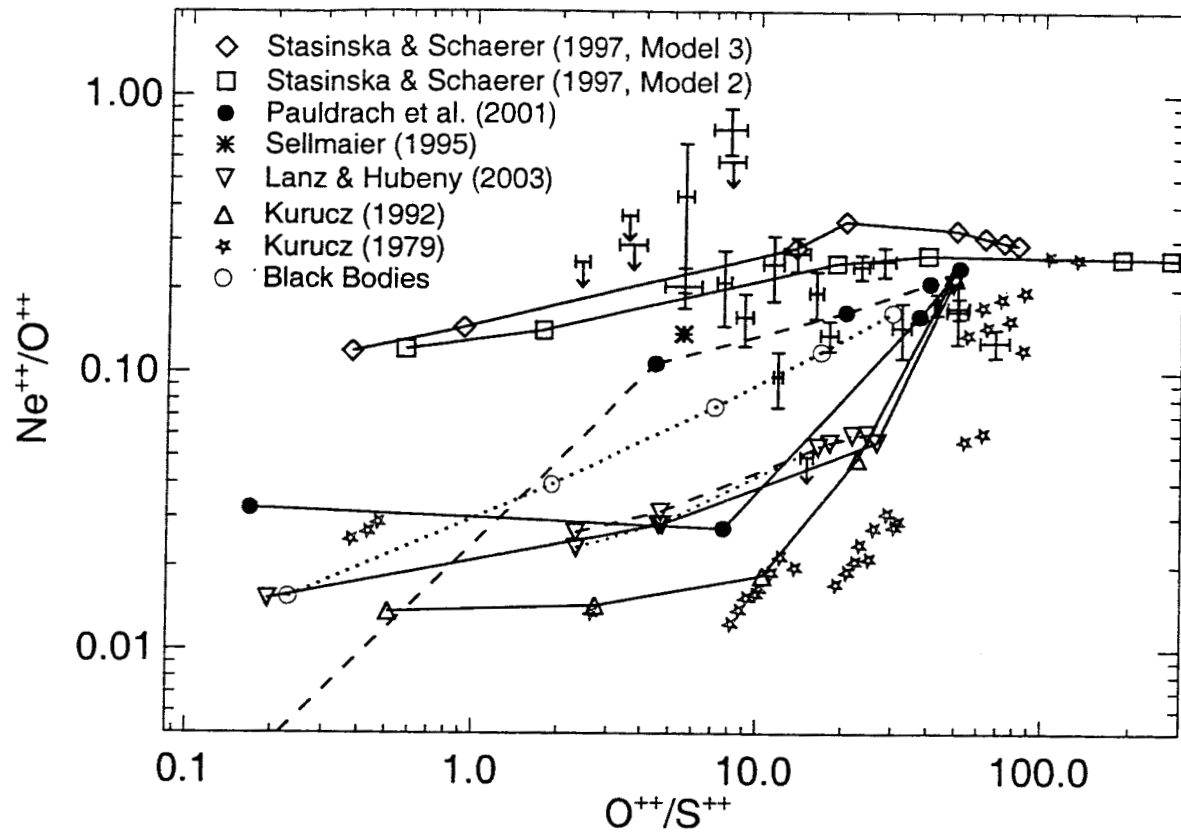


Table 1. Observed Line and Continuum Fluxes.

Line	Line Flux <sup>a</sup> ( $10^{-18}$ W cm <sup>-2</sup> ) Continuum Flux <sup>a</sup> (Jy)						
	G333.6-0.2				W43 (G30.8-0.0)		
	G333.6-ISO <sup>b</sup>	G333.6-N1 <sup>c</sup>	G333.6-N2 <sup>c</sup>	G333.6-N3 <sup>c</sup>	W43N <sup>c</sup>	W43C <sup>c</sup>	W43S <sup>c</sup>
[S III] 18.7 $\mu$ m	...	160.1 $\pm$ 16.2	36.3 $\pm$ 6.4	...	SCREH	SCREH	SCREH
	...	1370 $\pm$ 275	290 $\pm$ 280	...	SCREH	SCREH	SCREH
[S III] 33.5 $\mu$ m	...	189.1 $\pm$ 3.6	65.7 $\pm$ 4.7	38.9 $\pm$ 3.1	SCREH	SCREH	SCREH
	...	7040 $\pm$ 580	740 $\pm$ 185	730 $\pm$ 180	SCREH	SCREH	SCREH
[O III] 51.8 $\mu$ m	178.2 $\pm$ 8.5	84.3 $\pm$ 2.6	13.1 $\pm$ 1.6	6.8 $\pm$ 1.2	SCREH	SCREH	SCREH
	54400	9910 $\pm$ 320	1130 $\pm$ 360	410 $\pm$ 270	SCREH	SCREH	SCREH
[N III] 57.3 $\mu$ m	46.3 $\pm$ 4.9	51.1 $\pm$ 2.8	11.9 $\pm$ 1.4	...	SCREH	SCREH	SCREH
	57630	10945 $\pm$ 230	1310 $\pm$ 200	...	SCREH	SCREH	SCREH
[O III] 88.4 $\mu$ m	39.9 $\pm$ 3.9	37.7 $\pm$ 0.8	14.5 $\pm$ 0.9	7.5 $\pm$ 0.7	SCREH	SCREH	SCREH
	55100	9030 $\pm$ 175	1060 $\pm$ 270	630 $\pm$ 90	SCREH	SCREH	SCREH
[N II] 121.9 $\mu$ m	...	7.4 $\pm$ 1.4	2.6 $\pm$ 0.8	...	14.4 $\pm$ 0.8	5.5 $\pm$ 0.5	7.9 $\pm$ 0.8
	...	6790 $\pm$ 120	575 $\pm$ 230	...	5270 $\pm$ 420	1450 $\pm$ 260	4690 $\pm$ 380
[N II] 121.9 $\mu$ m <sup>d</sup>	...	...	...	...	4.9 $\pm$ 0.5 <sup>d</sup>	2.5 $\pm$ 0.4 <sup>d</sup>	2.6 $\pm$ 0.6 <sup>d</sup>
	...	...	...	...	2530 $\pm$ 210 <sup>d</sup>	800 $\pm$ 200 <sup>d</sup>	1600 $\pm$ 300 <sup>d</sup>
[N II] 205.2 $\mu$ m	...	2.0 $\pm$ 0.2	1.5 $\pm$ 0.2	1.7 $\pm$ 0.2	1.52 $\pm$ 0.26	0.46 $\pm$ 0.14	0.56 $\pm$ 0.13
	...	2000 $\pm$ 145	230 $\pm$ 120	115 $\pm$ 100	1300 $\pm$ 70	490 $\pm$ 50	1440 $\pm$ 40

<sup>a</sup>Flux and continuum errors are 1  $\sigma$  and are statistical only.

<sup>b</sup>All ISO LWS data had entrance aperture 80".

<sup>c</sup>All CGS data had entrance aperture 60" except as marked.

<sup>d</sup>Entrance aperture 45".



Table 2. Electron Densities and Abundance Ratios

	G333.6-0.2					W43		
	G333.6-ISO	G333.6-C <sup>a</sup>	G333.6-N1	G333.6-N2	G333.6-N3	W43N	W43C	W43S
$\tau_{9.6}$	1.5	1.5	1.5	1.5	1.5	2.74	2.74	3.63
$S_{5\text{GHz}}^b$	40.75	26.65	5.35	1.10	0.35	5.91	2.25	3.11
Electron Densities ( $\text{cm}^{-3}$ )								
$N_e$ (S III)	...	$2400 \pm 350^c$	$735 \pm 155^d$	$335 \pm 150^d$	...	$340 \pm 120^c$	$380 \pm 130^c$	$280 \pm 160^c$
$N_e$ (O III)	$2520 \pm 540^e$	$4200 \pm 700^c$	$730 \pm 50^d$	$140 \pm 50^d$	$140 \pm 70^d$	$870 \pm 50^c$	$730 \pm 60^c$	$725 \pm 50^c$
$N_e$ (N II)	...	$330^{+475}_{-185}^c$	$140^{+65}_{-50}^d$	$37^{+24}_{-21}^d$	...	$5120^{+\infty}_{-4200}^d$	$\infty^d$	$\infty^d$
$N_e$ (N II)	...	...	...	...	...	$180 \pm 80^f$	$650^{+7940}_{-440}^f$	$420^{+835}_{-250}^f$
Ionic Abundance Ratios ( $\times 10^{-6}$ )								
$\text{S}^{++}/\text{H}^+$	...	$4.3 \pm 0.7$	$12.2 \pm 1.6$	$14.2 \pm 2.1$	$20.8 \pm 2.7$	$9.0 \pm 1.4$	$14.1 \pm 2.2$	$12.8 \pm 2.3$
$\text{O}^{++}/\text{H}^+$	$38 \pm 5$	$71 \pm 9$	$80 \pm 8$	$50 \pm 6$	$82 \pm 12$	$130 \pm 14$	$275 \pm 31$	$227 \pm 25$
$\text{N}^{++}/\text{H}^+$	$14 \pm 2$	$36 \pm 4$	$48 \pm 5$	$27 \pm 4$	...	$68 \pm 8$	$127 \pm 14$	$117 \pm 13$
Ionic Abundance Ratios								
$\text{N}^{++}/\text{O}^{++}$	$0.36 \pm 0.04$	$0.51 \pm 0.02$	$0.60 \pm 0.04$	$0.53 \pm 0.07$	...	$0.52 \pm 0.03$	$0.46 \pm 0.02$	$0.52 \pm 0.02$
$\text{O}^{++}/\text{S}^{++}$	...	$16.7 \pm 2.2$	$6.53 \pm 0.55$	$3.56 \pm 0.57$	$3.95 \pm 0.52$	$14.5 \pm 1.2$	$19.5 \pm 1.9$	$17.8 \pm 1.8$
$\text{N}^{++}/\text{S}^{++}$	...	$8.47 \pm 0.70$	$3.90 \pm 0.29$	$1.89 \pm 0.29$	...	$7.54 \pm 0.38$	$9.04 \pm 0.36$	$9.18 \pm 0.37$
$\text{N}^+/\text{S}^{++}$	...	$0.63 \pm 0.20$	$1.90 \pm 0.52$	$1.75 \pm 0.45$	$14.3 \pm 1.3$	$1.83 \pm 0.37^g$	$3.98 \pm 1.67^g$	$2.30 \pm 1.38^g$
$\text{N}^{++}/\text{N}^+$	...	$13.3 \pm 4.0$	$2.06 \pm 0.45$	$1.08 \pm 0.27$	...	$4.11 \pm 0.72^g$	$2.27 \pm 0.84^g$	$3.99 \pm 1.46^g$
Column Density Ratios								
$\text{O}^{++}/\text{S}^{++}$	...	$9.55 \pm 0.43$	$6.56 \pm 0.22$	$8.60 \pm 0.87$	$3.95 \pm 0.52$	$5.64 \pm 0.22$	$10.2 \pm 0.5$	$6.89 \pm 0.47$
$\text{N}^{++}/\text{S}^{++}$	...	$4.85 \pm 0.23$	$3.92 \pm 0.24$	$4.56 \pm 0.67$	...	$2.94 \pm 0.15$	$4.72 \pm 0.18$	$3.56 \pm 0.19$
$\text{N}^+/\text{S}^{++}$	...	$4.62 \pm 0.58$	$9.96 \pm 1.29$	$15.7 \pm 2.4$	$14.3 \pm 1.3$	$4.61 \pm 0.36^h$	$3.29 \pm 0.37^h$	$1.95 \pm 0.50^h$
$\text{N}^{++}/\text{N}^+$	...	$1.05 \pm 0.13$	$0.39 \pm 0.08$	$0.29 \pm 0.09$	...	$0.64 \pm 0.05^h$	$1.44 \pm 0.14^h$	$1.27 \pm 0.14^h$
$\text{N}^+/\text{S}^{++}$	...	...	...	...	...	$3.39 \pm 0.41^g$	$2.31 \pm 0.37^g$	$1.55 \pm 0.34^g$
$\text{N}^{++}/\text{N}^+$	...	...	...	...	...	$0.87 \pm 0.12^g$	$2.04 \pm 0.34^g$	$2.30 \pm 0.52^g$
Adopted Abundance Ratios								
$\text{O}^{++}/\text{S}^{++}$	...	$13.1 \pm 3.6^i$	$6.53 \pm 0.55$	$3.56 \pm 0.57$	$3.95 \pm 0.52$	$14.5 \pm 1.2$	$19.5 \pm 1.9$	$17.8 \pm 1.8$
$\text{N}^{++}/\text{S}^{++}$	...	$6.66 \pm 1.81^i$	$3.90 \pm 0.29$	$1.89 \pm 0.29$	...	$7.54 \pm 0.38$	$9.04 \pm 0.36$	$9.18 \pm 0.37$
$\text{N}^+/\text{S}^{++}$	...	$2.63 \pm 1.99^i$	$5.93 \pm 4.03^i$	$8.71 \pm 6.96^i$	$14.3 \pm 6.0^i$	$2.96 \pm 0.38^j$	$3.02 \pm 0.80^j$	$2.59 \pm 0.74^j$
$\text{N}^{++}/\text{N}^+$	...	$1.05 \pm 0.13^k$	$1.23 \pm 0.83^i$	$0.69 \pm 0.39^i$	...	$1.43 \pm 0.30^i$	$1.44 \pm 0.33^i$	$1.89 \pm 0.53^i$
$(\text{N}^+ + \text{N}^{++})/\text{S}^{++}$	...	$9.3 \pm 3.8$	$9.8 \pm 4.3$	$10.6 \pm 7.3$	$>14 \pm 6$	$10.5 \pm 0.8$	$12.1 \pm 1.2$	$11.8 \pm 1.1$

<sup>a</sup>Center position, data from Colgan et al. (1993).

<sup>b</sup>Radio flux in the aperture used for the FIR line measurement, scaled to the value appropriate for  $\nu = 5$  GHz.

<sup>c</sup>Electron density computed using the 45'' aperture data.

<sup>d</sup>Electron density computed using the 60'' aperture data.

<sup>e</sup>Electron density computed using the 80'' aperture data.

<sup>f</sup>Electron density computed using the 45'' aperture data for the 122  $\mu\text{m}$  line and reducing the 205  $\mu\text{m}$  line flux by the ratio of the aperture diameters.

<sup>g</sup>Computed using the 45'' aperture data for the 122  $\mu\text{m}$  line and reducing the 205  $\mu\text{m}$  line flux by the ratio of the aperture diameters.

<sup>h</sup>Computed using the 60'' aperture data for the 122  $\mu\text{m}$  line and reducing both the 122  $\mu\text{m}$  and the 205  $\mu\text{m}$  line fluxes by the ratio of the aperture diameters for comparison to the [N III] and [S III] lines.

<sup>i</sup>Average of column density ratio and ionic abundance ratio; the error is increased for G333.6-N3 where  $N_e(\text{N II})$  was not measured.

<sup>j</sup>Average of ionic abundance ratios for all 122  $\mu\text{m}$  line measurements plus column density for the 45'' aperture 122  $\mu\text{m}$  line measurements.

Table 3. Estimated Abundances

Element	G333.6-0.2	W43	Orion Nebula
Estimated S/H <sup>a</sup>	$13.4 \pm 0.8 \times 10^{-6}$	$12.1 \pm 1.1 \times 10^{-6}$	$7.0 \times 10^{-6}$ <sup>b</sup>
Estimated N/S <sup>a</sup>	$9.3 \pm 2.1$	$10.3 \pm 0.4$	$9.71$ <sup>b</sup>
Estimated N/H <sup>c</sup>	$12.5 \times 10^{-5}$	$12.5 \times 10^{-5}$	$6.8 \times 10^{-5}$ <sup>d</sup>
Estimated Ne/H	$3.4 \times 10^{-4}$ <sup>e</sup>	$1.1 \times 10^{-4}$ <sup>f</sup>	$1.0 \times 10^{-4}$ <sup>b</sup>
Estimated O/H	$3 \text{ to } 4 \times 10^{-4}$ <sup>g</sup>	$3 \text{ to } 8 \times 10^{-4}$ <sup>g</sup>	$4.0 \times 10^{-4}$ <sup>d</sup>

<sup>a</sup> Assumed fractional ionization  $\langle S^{++}/S \rangle = 0.9$ .

<sup>b</sup> Simpson et al. (1998) revised for new collisional excitation cross sections (see text).

<sup>c</sup>  $N/H = N/S \times S/H$ .

<sup>d</sup> Rubin et al. (1991).

<sup>e</sup> Estimated from [Ne II] 12.8  $\mu\text{m}$  measurements of Fujiyoshi et al. (1998) and radio continuum measurements of Fujiyoshi (1999).

<sup>f</sup> SCREH.

<sup>g</sup>  $O/H = O^{++}/H^{+} \times ICF_{\text{Model}}$ .

Table 4. G333.6-0.2 - Comparison of Models to Observations

Line Flux or Ratio	Observed Fluxes and Ratios <sup>a</sup>				Model for G333.6-0.2 Core <sup>a,b</sup>				Model for North Positions <sup>a,b</sup>		
	Center	N1	N2	N3	Center	N1	N2	N3	N1	N2	N3
[O III] F <sub>52</sub> $\mu\text{m}$	161 <sup>c</sup>	92	14	7	350	0.0005	0	0	117	1	0
[O III] F <sub>88</sub> $\mu\text{m}$	27 <sup>c</sup>	39	15	8	54	0.0001	0	0	67	1	0
[O III] F <sub>52</sub> /F <sub>88</sub>	6.06	2.37	0.96	0.96	6.50	4.12	...	...	1.75	1.61	...
[S III] F <sub>19</sub> $\mu\text{m}$	412 <sup>c</sup>	270	61	...	869	273	50	14	154	52	16
[S III] F <sub>33</sub> $\mu\text{m}$	158 <sup>c</sup>	228	79	47	351	243	72	27	192	85	33
[S III] F <sub>19</sub> /F <sub>33</sub>	2.61	1.19	0.78	...	2.48	1.13	0.70	0.52	0.80	0.61	0.48
[N II] F <sub>122</sub> $\mu\text{m}$	2.5 <sup>c</sup>	7.4	2.6	...	9.2	13.7	5.6	2.8	6.8	6.9	3.5
[N II] F <sub>205</sub> $\mu\text{m}$	0.44 <sup>c</sup>	1.97	1.46	1.66	1.5	2.4	1.4	0.9	1.9	2.0	1.4
[N II] F <sub>122</sub> /F <sub>205</sub>	5.7	3.7	1.8	...	6.2	5.8	4.1	3.0	3.7	3.4	2.5
[N III] F <sub>57</sub> $\mu\text{m}$	52 <sup>c</sup>	55	13	...	115	0.006	0	0	62	1.9	0
[N III] F <sub>57</sub> /[N II] F <sub>122</sub>	21.0	7.4	5.0	...	12.6	0.0004	0	0	9.1	0.28	0
[N III] F <sub>57</sub> /[O III] F <sub>52</sub>	0.32	0.60	0.89	...	0.33	12	...	...	0.53	1.49	...
[Ne III] F <sub>36</sub> $\mu\text{m}$ <sup>c</sup>	17	...	...	...	8.5	...	...	...	...	...	...
[Ne II] F <sub>12.8</sub> $\mu\text{m}$ <sup>d</sup>	1300	...	...	...	1793	...	...	...	...	...	...
[S IV] F <sub>10.5</sub> $\mu\text{m}$ <sup>e</sup>	100	...	...	...	95	...	...	...	...	...	...
[Ar III] F <sub>9.0</sub> $\mu\text{m}$ <sup>f</sup>	350	...	...	...	251	...	...	...	...	...	...
[Ar III] F <sub>9.0</sub> $\mu\text{m}$ <sup>g</sup>	83	...	...	...	150	...	...	...	...	...	...
[Ar II] F <sub>7.0</sub> $\mu\text{m}$ <sup>g</sup>	74	...	...	...	13	...	...	...	...	...	...
He <sup>+</sup> /H <sup>+</sup> RRLs <sup>h</sup>	0.045	...	...	...	0.044	...	...	...	...	...	...
[O III] F <sub>52</sub> $\mu\text{m}$ <sup>i</sup>	194	...	...	...	359	...	...	...	...	...	...
[O III] F <sub>88</sub> $\mu\text{m}$ <sup>i</sup>	41	...	...	...	56	...	...	...	...	...	...
[N III] F <sub>57</sub> $\mu\text{m}$ <sup>i</sup>	50	...	...	...	144	...	...	...	...	...	...

<sup>a</sup>Fluxes  $\times 10^{-18}$  W cm<sup>-2</sup>, corrected for extinction.

<sup>b</sup>Model Abundances ( $\times 10^{-6}$ ): C/H = 400, N/H = 125, O/H = 500, Ne/H = 340, S/H = 13.4, Ar/H = 4.5.

<sup>c</sup>0.75' Aperture, Colgan et al. (1993).

<sup>d</sup>0.75' Aperture, Geballe et al. (1981).

<sup>e</sup>0.50' Aperture, Geballe et al. (1981).

<sup>f</sup>0.67' Aperture, Geballe et al. (1981).

<sup>g</sup>0.35' Aperture, Cohen et al. (1989), Simpson et al. (1995b).

<sup>h</sup>Radio recombination lines measured in 2.3' aperture, McGee & Newton (1981).

<sup>i</sup>80'' Aperture for ISO LWS.

Table 5. W43 - Comparison of Models to Observations

Ratio	Observed Ratios	Model D <sup>a</sup>	Model S <sup>a</sup>	Model T <sup>a</sup>	Comments
He <sup>+</sup> /H <sup>+</sup> RRLs	0.076 <sup>b</sup>	0.075	0.073	0.075	Good $T_{\text{eff}}$
[S III] 19/33 $\mu\text{m}$	0.78, 0.82, 0.72	1.04	1.03	1.02	Model $N_e$ too high
[O III] 52/88 $\mu\text{m}$	2.6, 2.3, 2.3	2.30	2.33	2.33	Good $N_e$
[N II] 122/205 $\mu\text{m}$	4.3, 7.1, 6.2	6.29	6.28	6.19	Good $N_e$
S <sup>++</sup> /S	...	0.919	0.953	0.965	Implies S/H = 11.8, 11.4, or $11.3 \times 10^{-6}$
O <sup>++</sup> /S <sup>++</sup> <sup>c</sup>	$16.33 \pm 0.89$	0.583	0.391	0.221	Implies O/H = 3.3, 4.8, or $8.3 \times 10^{-4}$ <sup>d</sup>
N <sup>++</sup> /O <sup>++</sup>	$0.50 \pm 0.01$	1.18	1.58	3.05	Implies N/O = 0.42, 0.32, or 0.16 <sup>e</sup>
N <sup>++</sup> /N <sup>+</sup>	1.5 to 3.0	1.72	1.45	1.86	Compare uniform He <sup>+</sup> /H <sup>+</sup>
N <sup>+</sup> /S <sup>++</sup>	$2.90 \pm 0.31$	0.400	0.428	0.363	Implies N/S = 7.25, 6.78, or 7.99 <sup>f</sup>
N <sup>++</sup> /S <sup>++</sup>	$8.61 \pm 0.21$	0.686	0.620	0.673	Implies N/S = 12.55, 13.89, or 12.79 <sup>f</sup>

<sup>a</sup>The stellar atmosphere used for the Model D is the sum of 0.8 times the D-35 model atmosphere from Pauldrach et al. (2001) plus 0.2 times the D-40 model atmosphere, the stellar atmosphere used for Model S is the sum of 0.6 times the S-35 model atmosphere plus 0.4 times the S-40 model atmosphere, where D-35, D-40, S-35, and S-40 stand for Dwarf and Supergiant models with  $T_{\text{eff}}$  equal to 35,000 K and 40,000 K, respectively, and the stellar atmosphere used for Model T is the TLUSTY model atmosphere from Lanz & Hubeny (2003) for  $T_{\text{eff}} = 35,000$  K and  $\log g = 3.50$ . The model abundance ratios ( $\times 10^{-6}$ ) are C/H = 400, N/H = 125, O/H = 500, Ne/H = 110, S/H = 12.1, and Ar/H = 4.5.

<sup>b</sup>Weighted average of measurements of Churchwell et al. (1978), Lichten, Rodriguez, & Chaisson (1979), Lockman & Brown (1982), Peimbert et al. (1992).

<sup>c</sup>Model O<sup>++</sup>/S<sup>++</sup> is the density weighted ratio of  $\langle \text{O}^{++}/\text{O} \rangle / \langle \text{S}^{++}/\text{S} \rangle$ , etc. (SCREH). It should be multiplied by the relevant abundance ratio for comparison with the observed ratios.

<sup>d</sup>O/H = O/S  $\times$  S/H.

<sup>e</sup>N/H = 1.4, 1.5, or  $1.4 \times 10^{-4}$  if O/H = 3.3, 4.8, or  $8.3 \times 10^{-4}$ .

<sup>f</sup>The average N/S = 9.9, 10.3, or 10.4, N/H = 1.17, 1.18, or  $1.17 \times 10^{-4}$ , and O/H = 2.8, 3.7, or  $7.2 \times 10^{-4}$ .